

Measurement of NC $e^\pm p$ cross sections at high Bjorken x with the ZEUS detector



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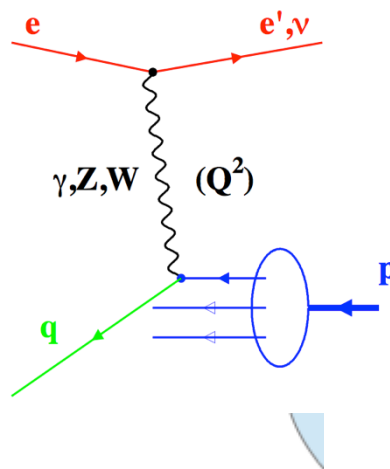


On behalf of the ZEUS collaboration

Resolving Structure of Matter

HERA, e (27.5 GeV)
p (920 GeV) collider
to study the proton
structure with a
high resolving
power. ($\sim 10^{-3}$ fm)

HERA NC(CC) Events



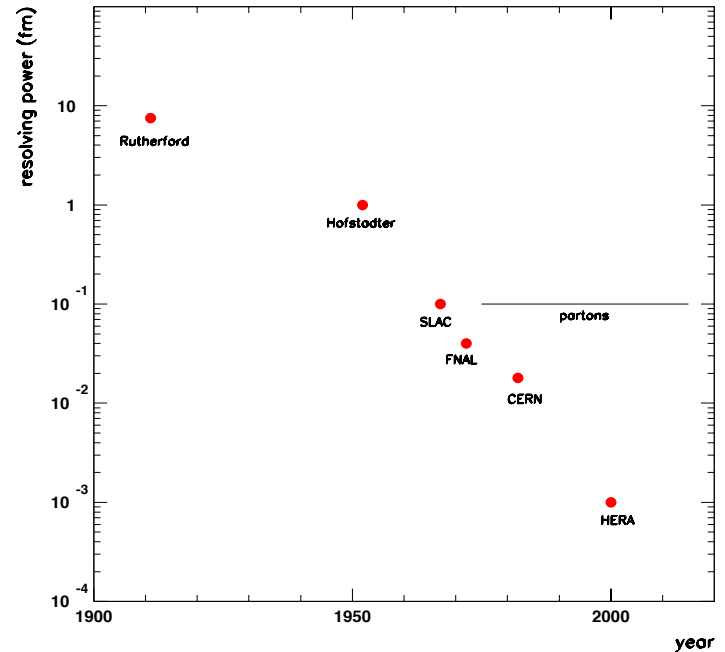
Kinematics:

$$Q^2 = -q^2 = -(k - k')^2$$

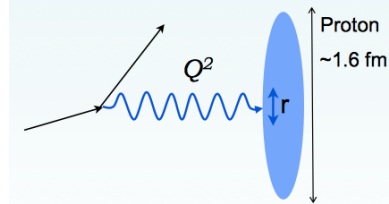
$$x = \frac{Q^2}{2p \cdot q}$$

$$y = \frac{p \cdot q}{p \cdot k}$$

Boson
virtuality
Bjorken
variable
Inelasticity



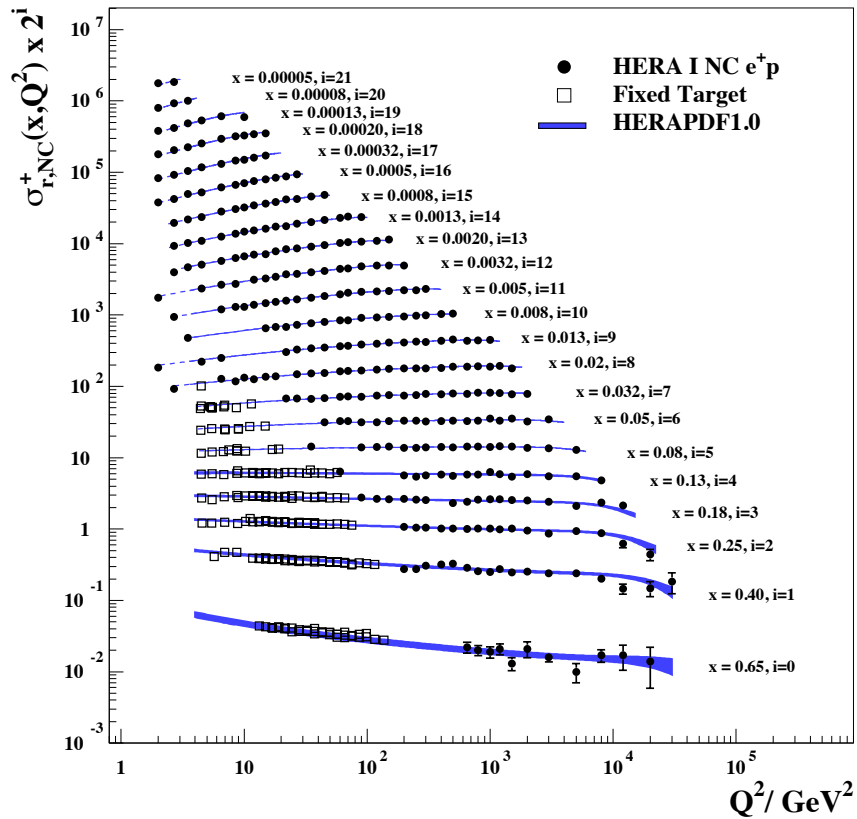
4-momentum transfer Q^2 defines
distance scale r at which proton is probed



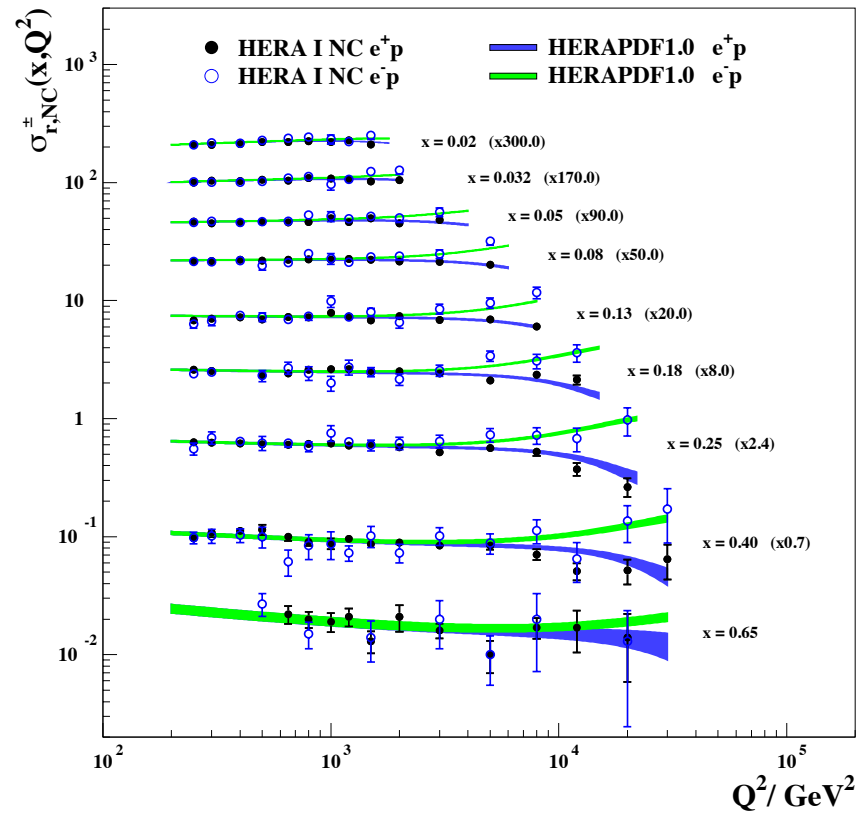
$$r \approx \hbar c / Q = 0.2[\text{fm}] / Q[\text{GeV}]$$

Combined cross sections, H1+ZEUS

H1 and ZEUS



H1 and ZEUS

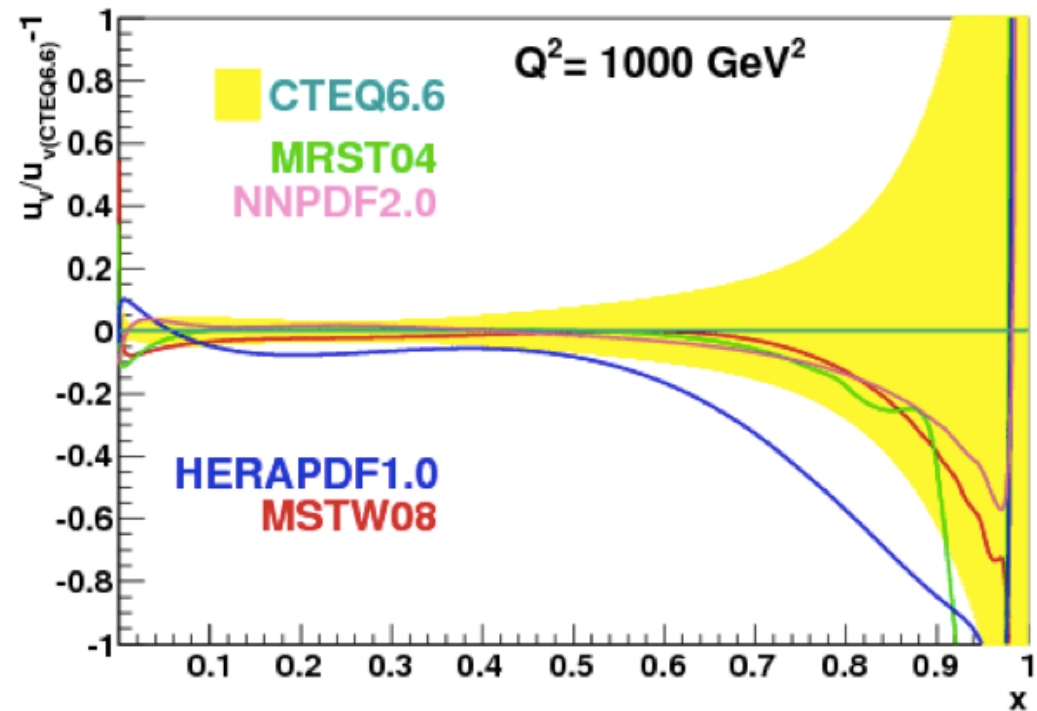


High x - Motivation

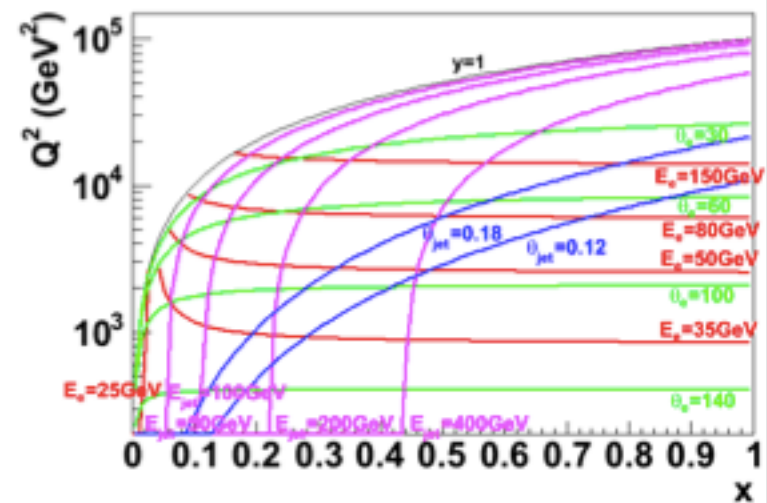
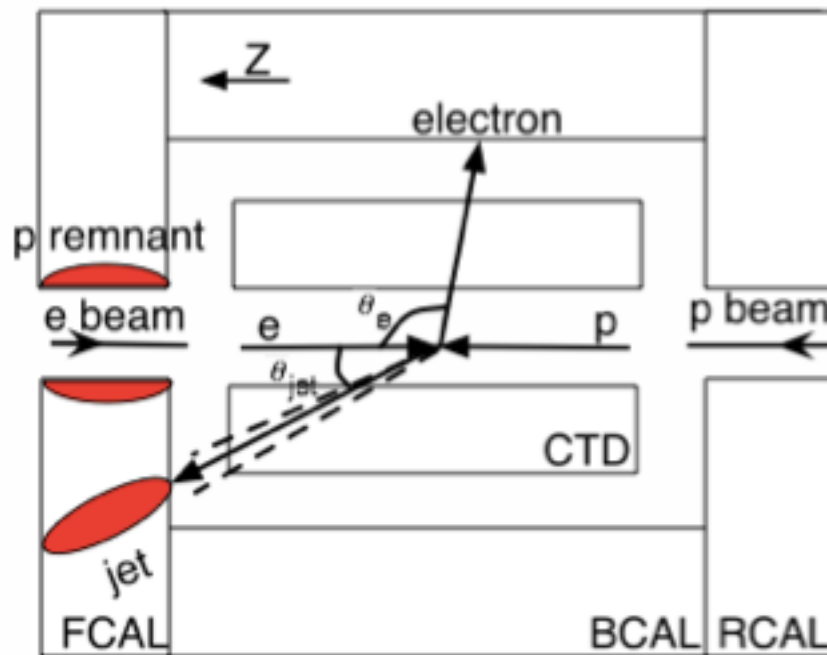
- At very low x , access to Q^2 limited.
- At low x , photon starts fluctuating long before resolving the proton. Study the γ^*p interaction region.
- As x increases, reach of higher Q^2 values is possible. Photon can resolve details of the proton.
- As x and Q^2 increase, probe's fluctuation time is short. Probe 'sees' inside the proton.
- HERA data, used for PDF determination, only up to $x=0.65$.
- Higher x data exist only from fixed target experiments and at low Q^2 .
- Want to measure at HERA at high Q^2 up to $x=1$ and use these data for constraining PDFs at high x .

High x - Motivation

The PDF's are poorly determined at high- x . Sizeable differences despite the fact that all fitters use the same parametrization $xq \propto (1-x)^\eta$. Is it possible to check this ?



HERA high x high Q²

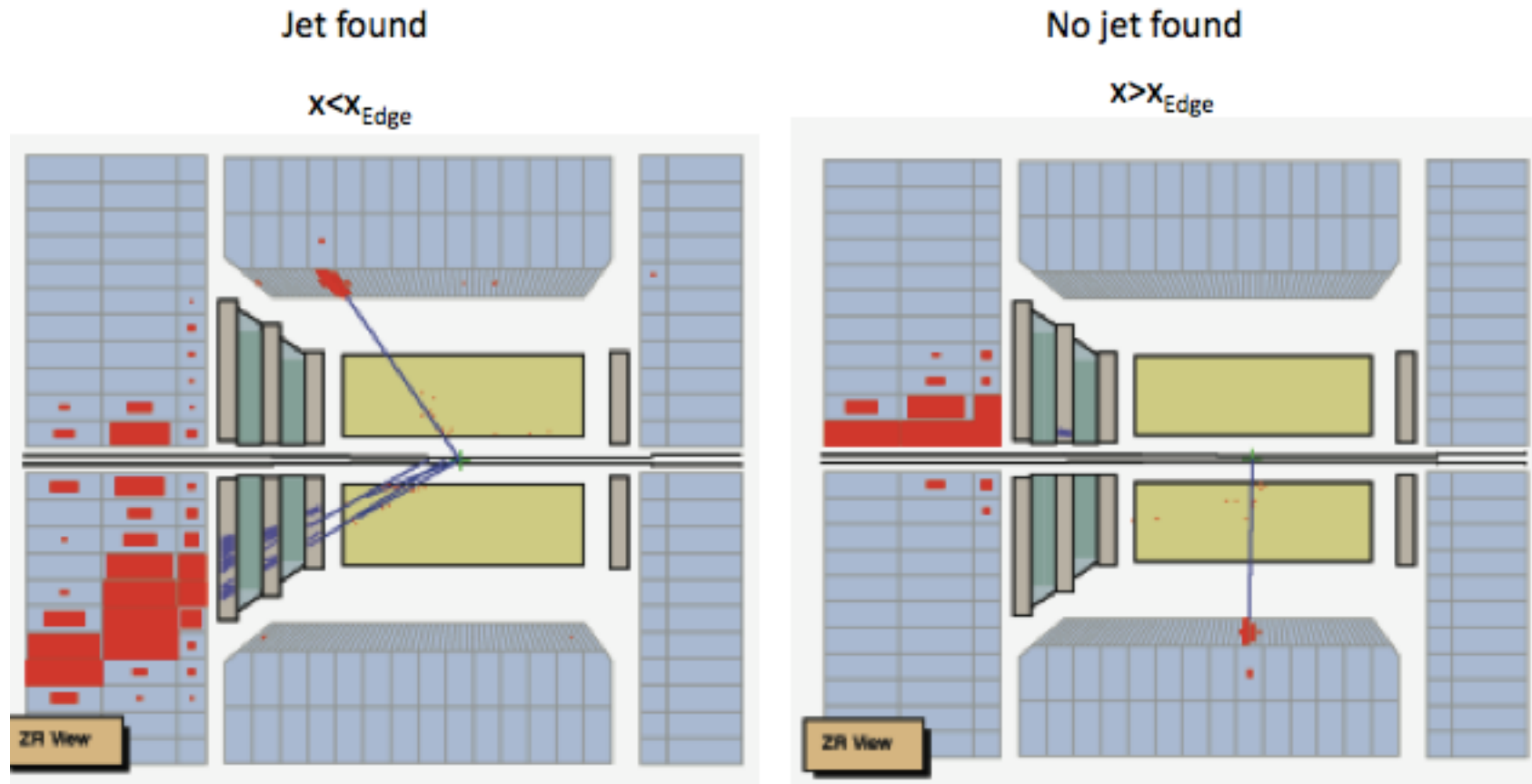


- At high Q², scattered electron seen with ≈100% acceptance

- For not too high x, measure x from jet: $\frac{d^2\sigma}{dx dQ^2}$

- For $x > x_{\text{Edge}}$, measure $\int_x \frac{d^2\sigma}{dx dQ^2} dx$

1-jet, 0-jet



x and Q² reconstruction

Electron Pt jet method is chosen for reconstruction of kinematic variables

$$Q_e^2 = 4E_e E'_e \cos^2 \frac{\theta_e}{2}, \quad Q^2 \text{ by electron method}$$

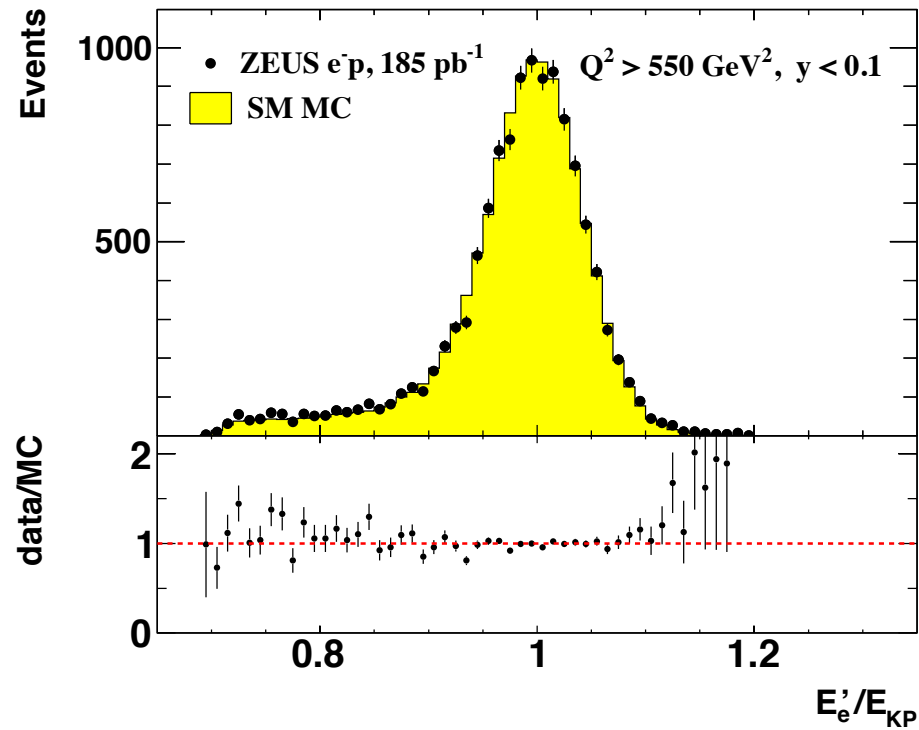
x for 1 jet events

$$x_{pt_e} = \frac{(p_{t_e}/\sin\theta_{jet})(1 + \cos\theta_{jet})}{2E_p \left(1 - \frac{(p_{t_e}/\sin\theta_{jet})(1 - \cos\theta_{jet})}{2E_e}\right)}$$

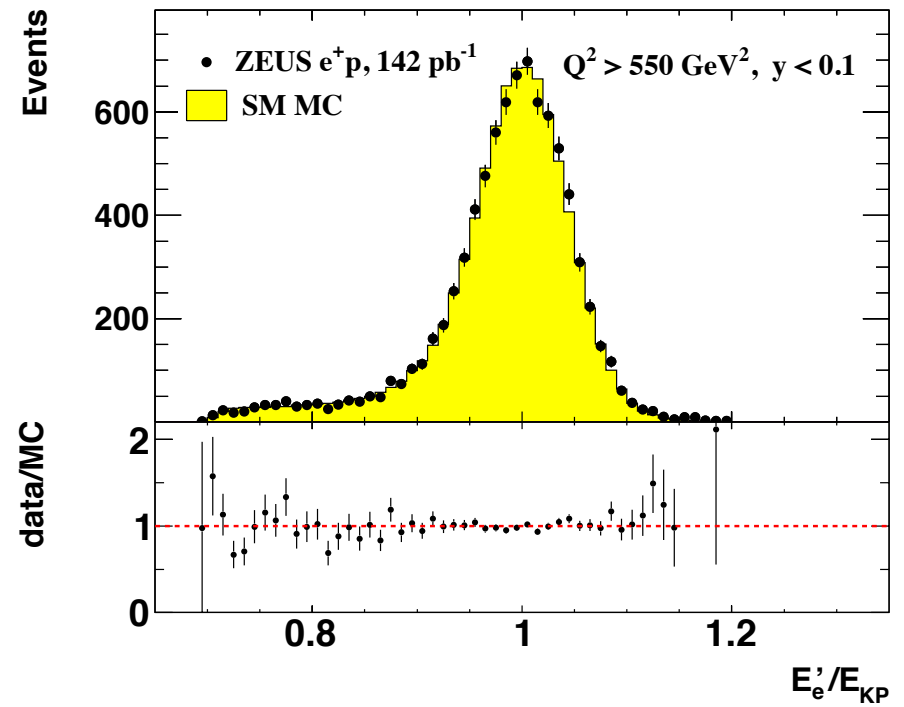
$$x = \frac{E_{jet}(1 + \cos\theta_{jet})}{2E_p \left(1 - \frac{E_{jet}(1 - \cos\theta_{jet})}{2E_e}\right)} \quad \text{x for multi-jet events}$$

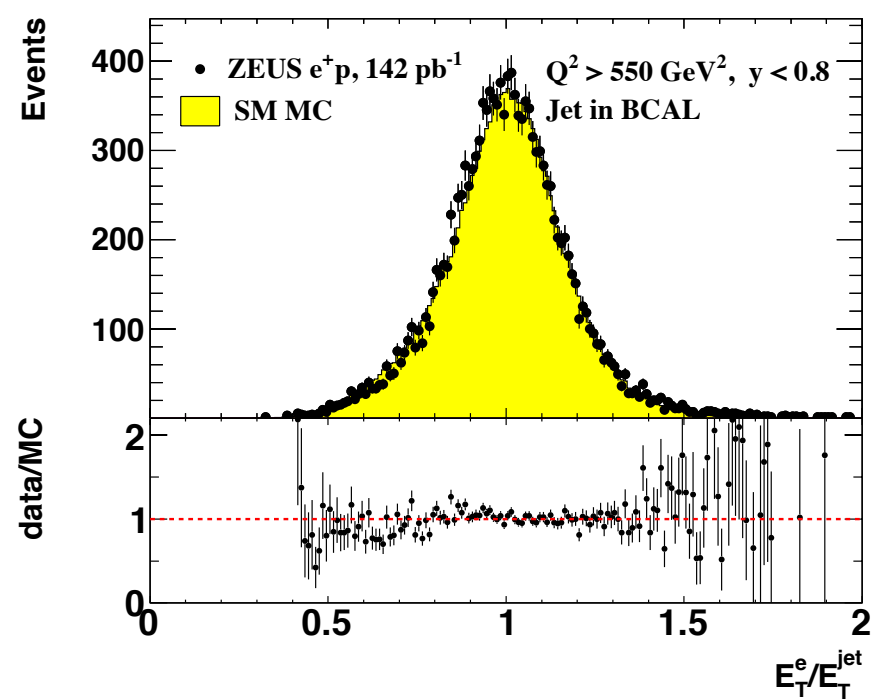
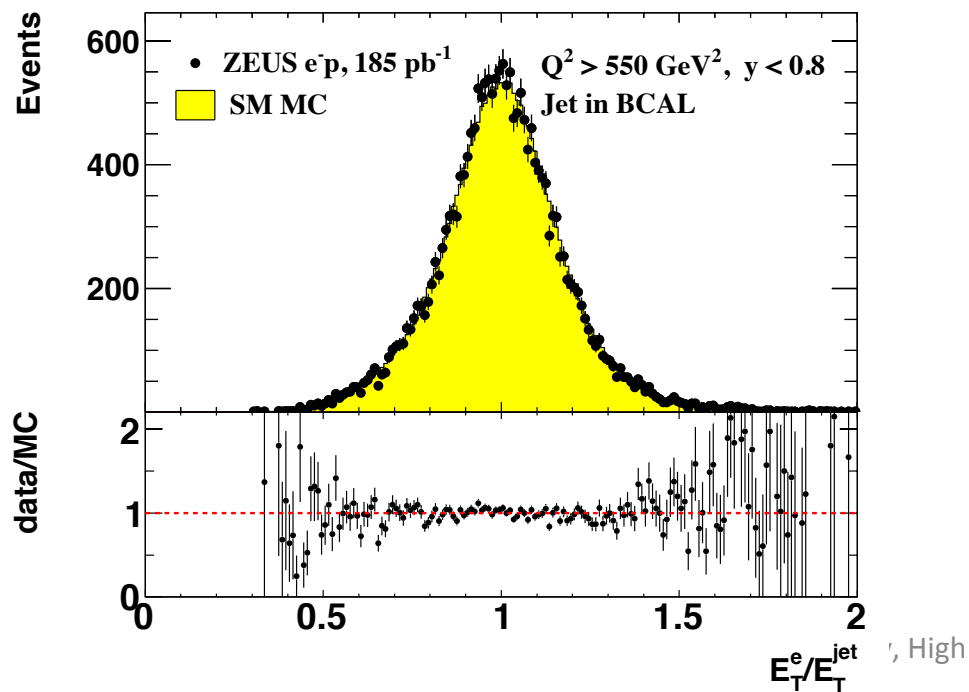
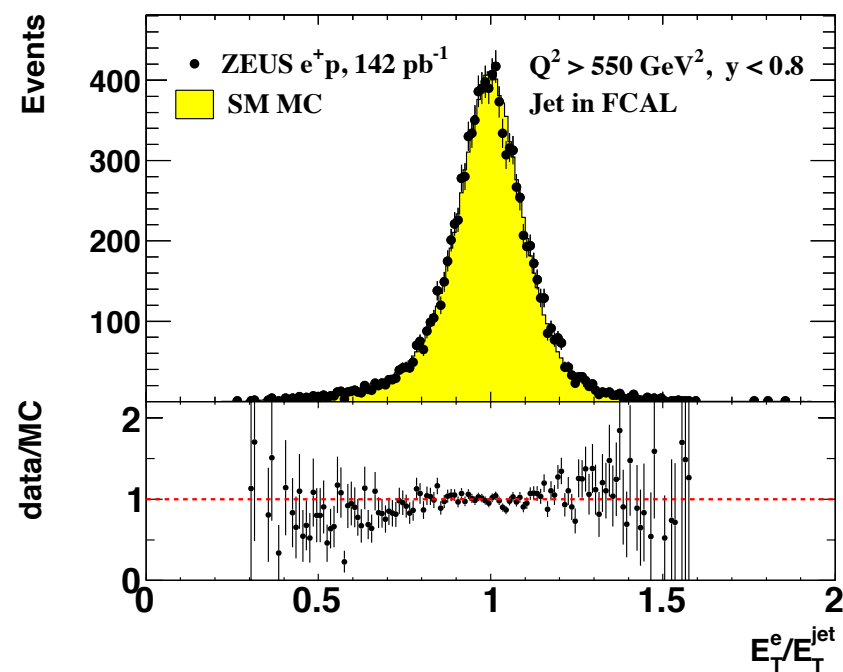
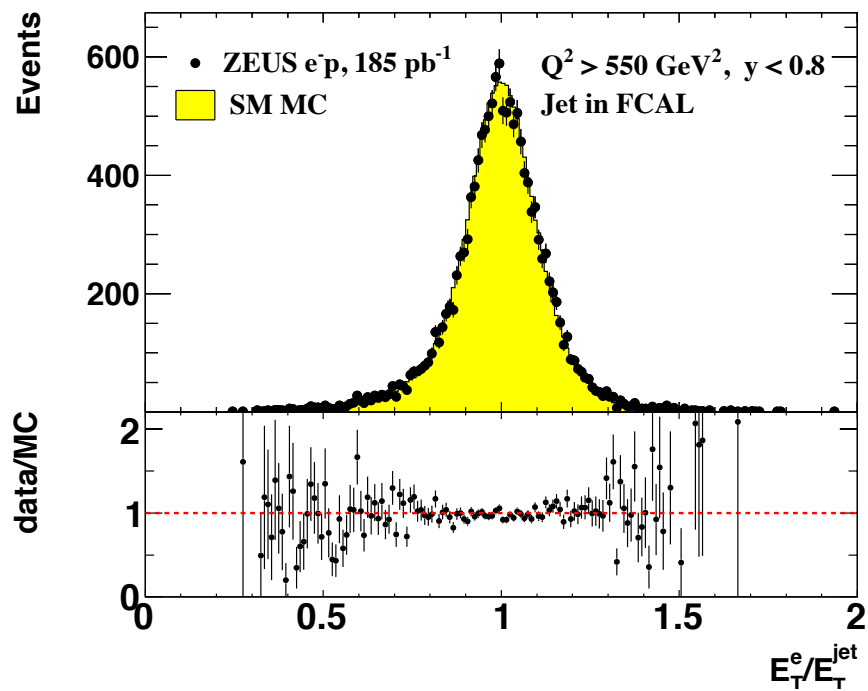
Control plots

e^-p



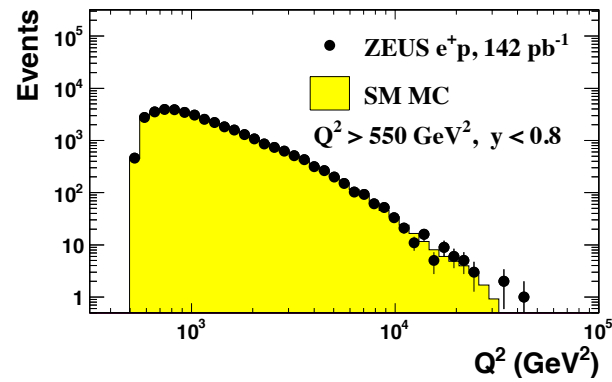
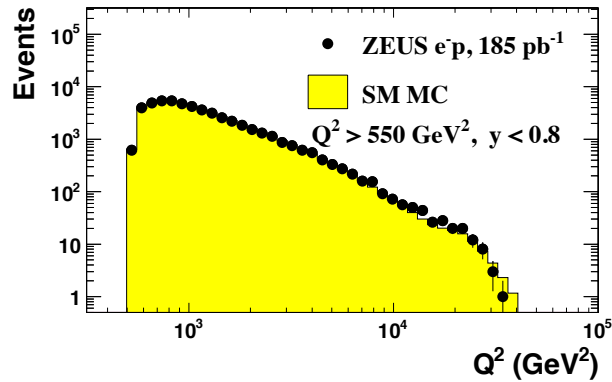
e^+p



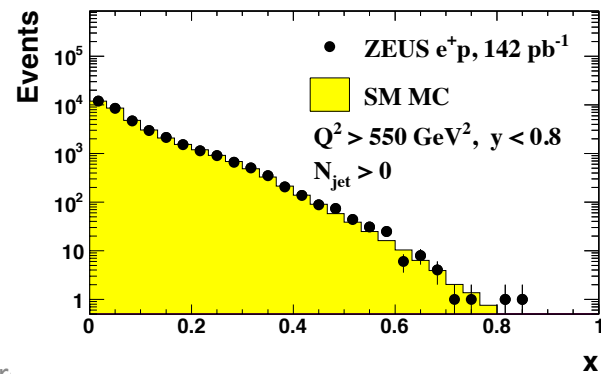
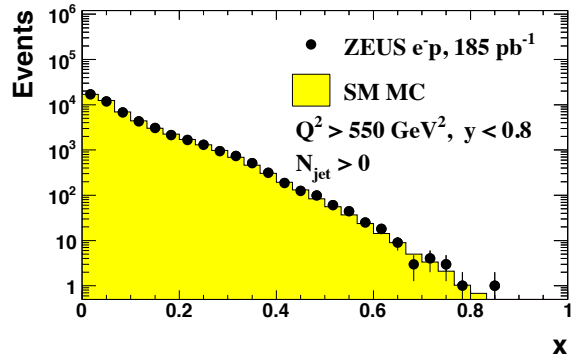
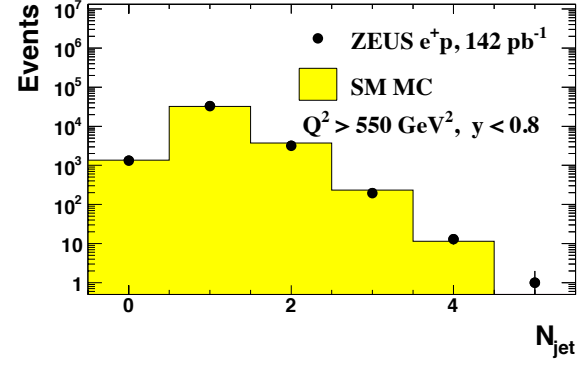
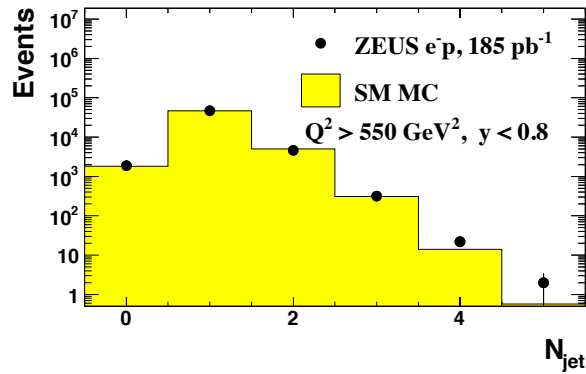


Control plots

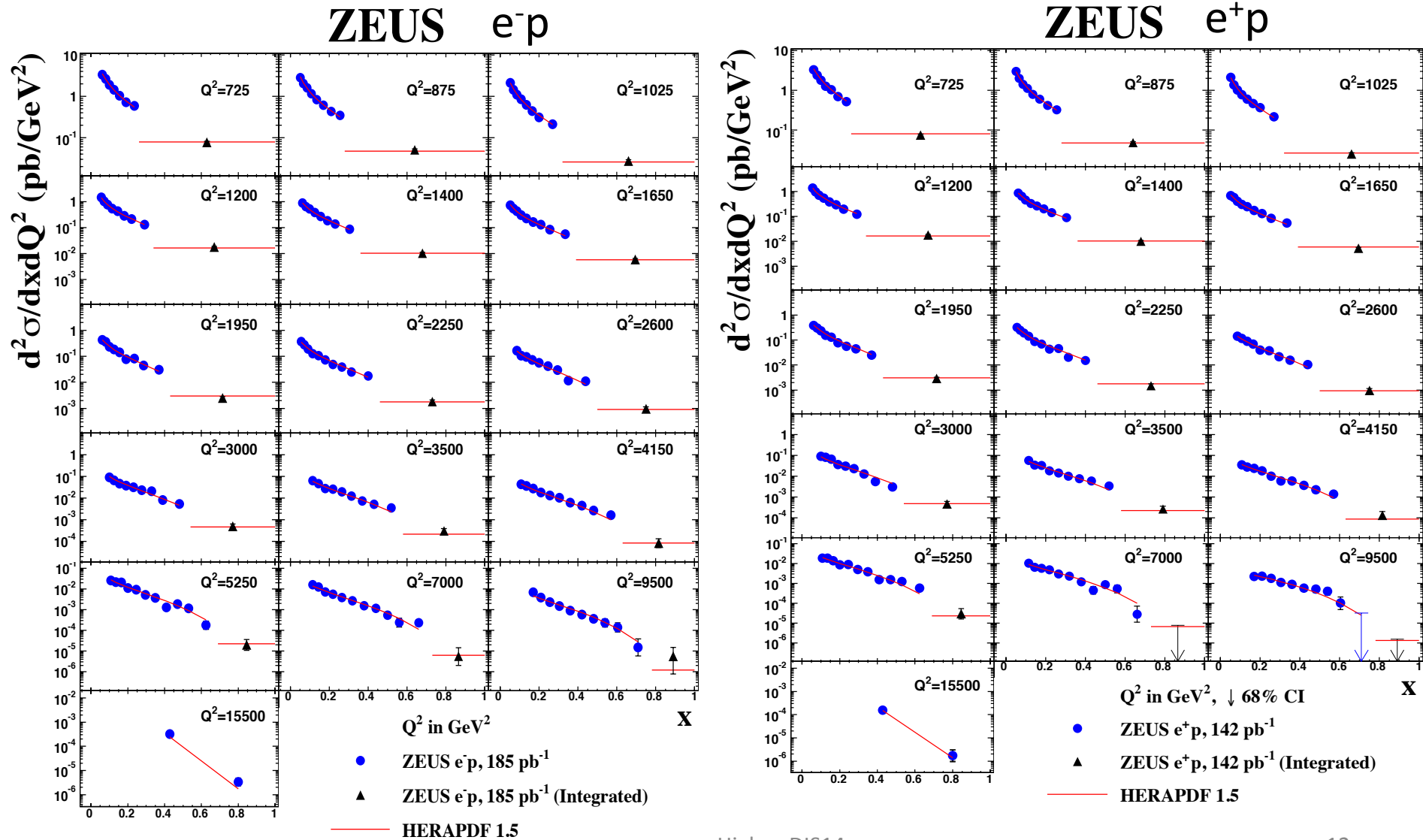
e^-p



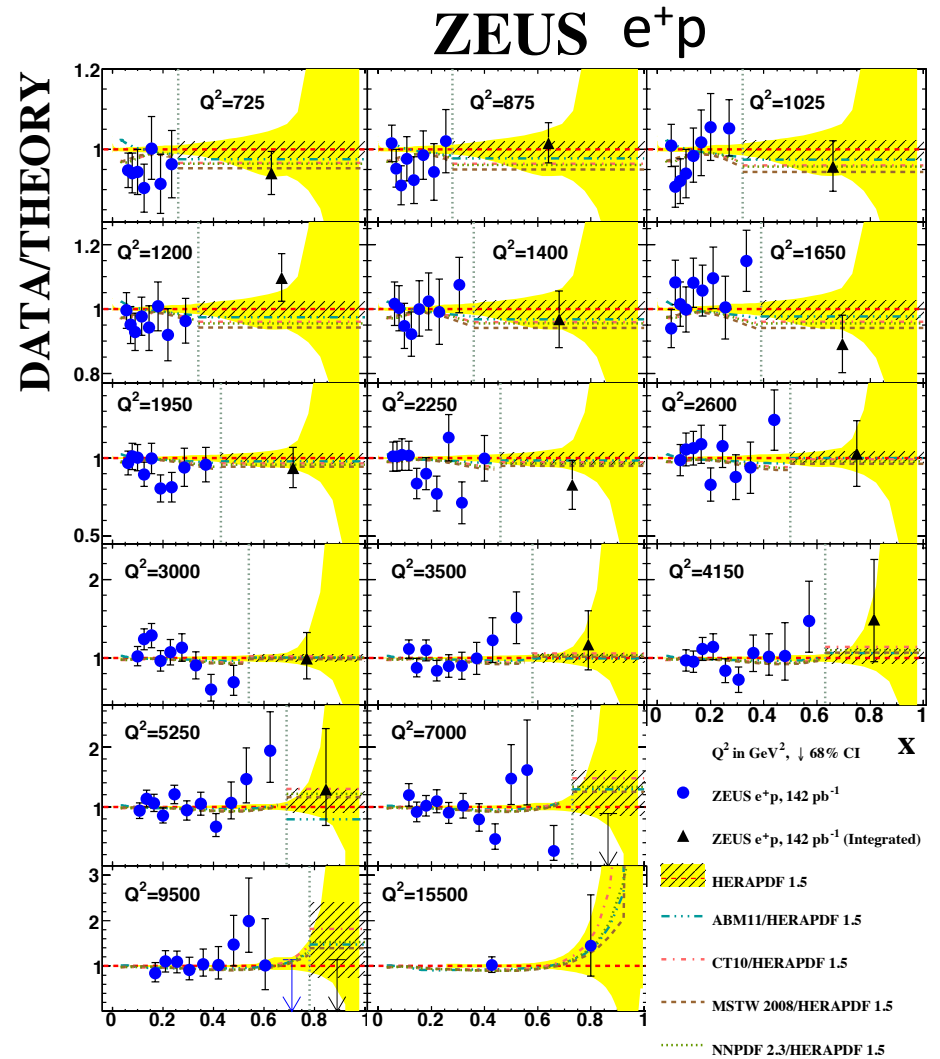
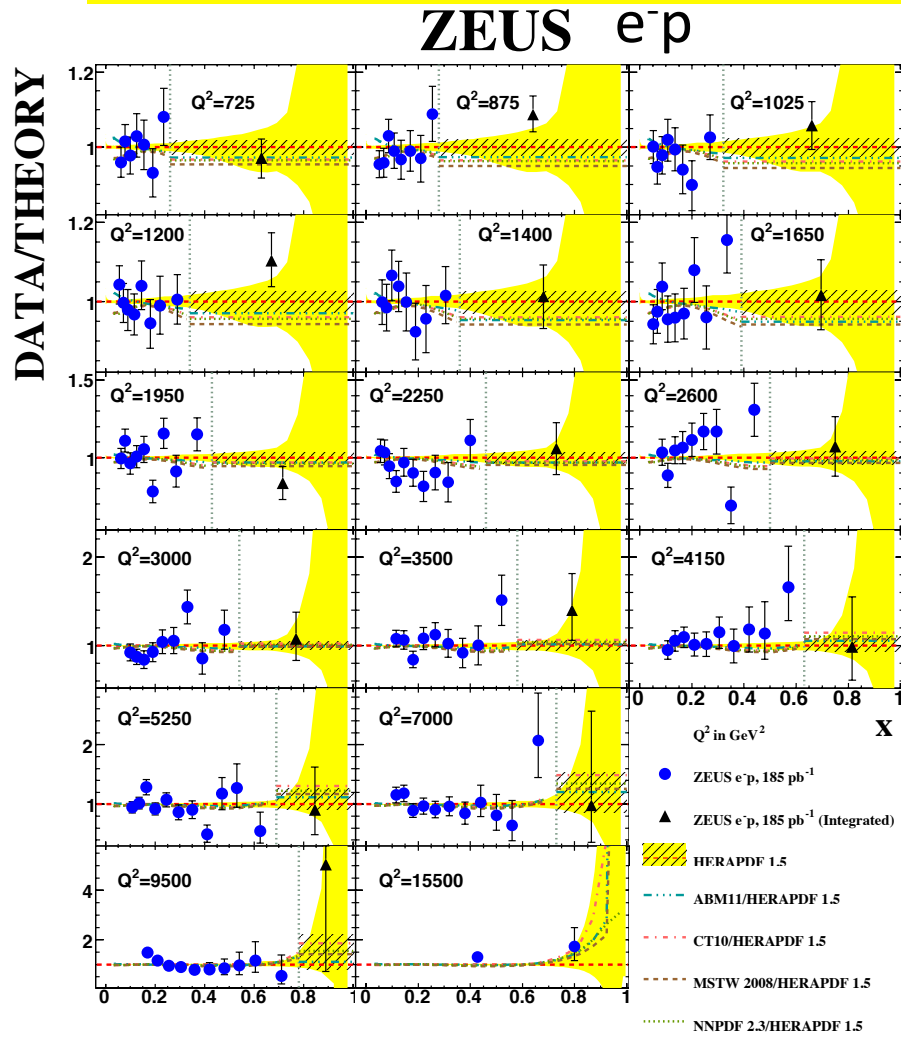
e^+p



NC cross section $e^\pm p$



Data/Theory NC $e^\pm p$



How to include these data for PDF fits?

Simple chi squared methods are not appropriate for including these data in fits. Many of the bins have only a few or zero events. One should use the number of measured events in the fits. Fitters should provide predictions for the number of events expected in the bins in which event counts are reported, and then use Poisson statistics to calculate the probability for the observed number of events given this expectation (from a Note by Allen Caldwell).

$Q^2(\text{GeV}^2)$	x	$N(e^-)$	$N(e^+)$	x_{edge}	$N(e^-)$	$N(e^+)$
3500	0.52	33	24	0.79	17	11
4150	0.57	18	12	0.81	5	6
5250	0.62	5	13	0.85	3	3
7000	0.66	10	1	0.87	1	0
9500	0.71	1	0	0.89	1	0
15500	0.80	8	3			

Conclusions high x

Measured $e^\pm p$ NC DIS cross sections at $Q^2 > 725 \text{ GeV}^2$ up to $x \approx 1$.

Fine binning in x , extension of kinematic coverage up to $x \approx 1$ make data important input to fits constraining the PDFs in the valence-quark domain.

More details: Phys. Rev. D89 (2014) 072007
([arXiv:1312.4438 \[hep-ex\]](https://arxiv.org/abs/1312.4438))