

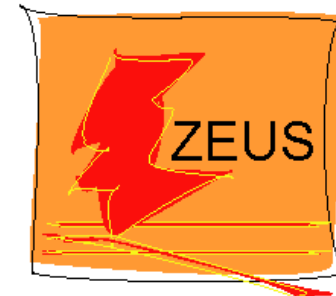
Jet and photon production and extraction of α_S at HERA

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On behalf of H1 and ZEUS collaborations

¹ Deutsches Elektronen-Synchrotron (DESY)

New Frontiers in Physics
Kolymbari, August 24



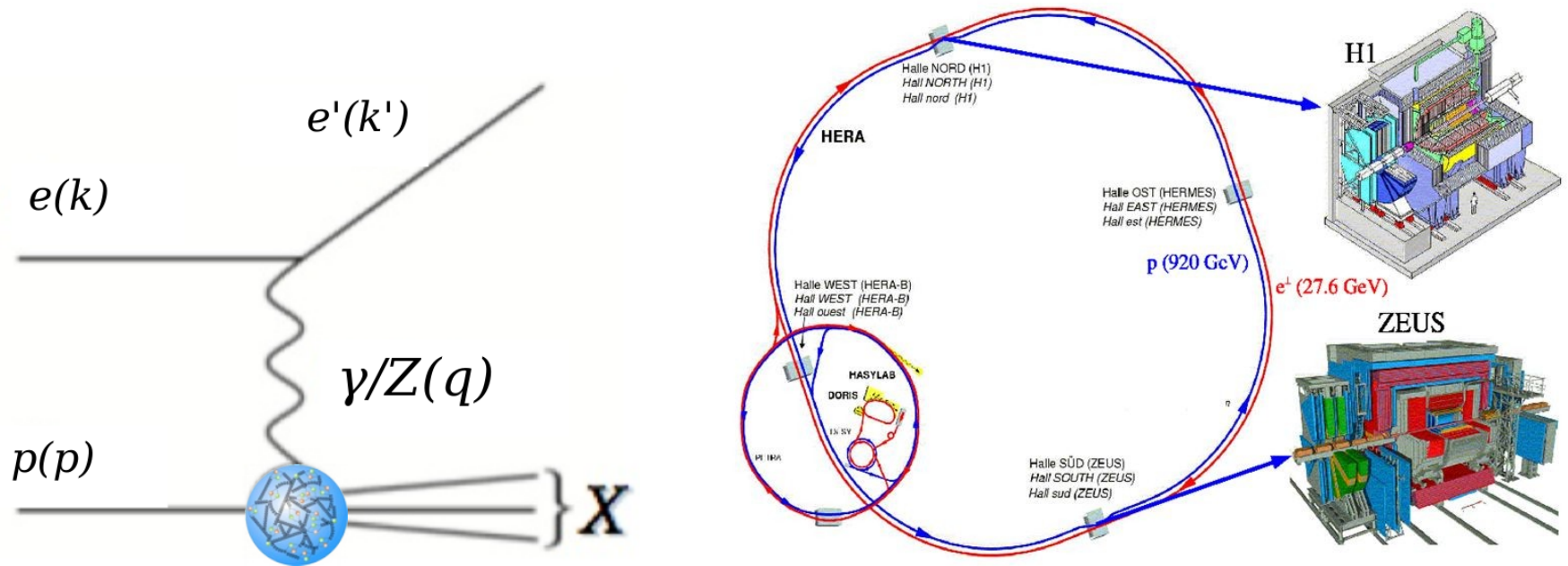
HERA Collider

- The only existing ep collider (1992 - 2007)
- About **0.5 fb⁻¹** of data per experiment
- Two multi-purpose detectors (**H1 + ZEUS**)

$$e^\pm + p$$

$$27.6 \text{ GeV} + 920 \text{ GeV}$$

$$\sqrt{s} = 319 \text{ GeV}$$

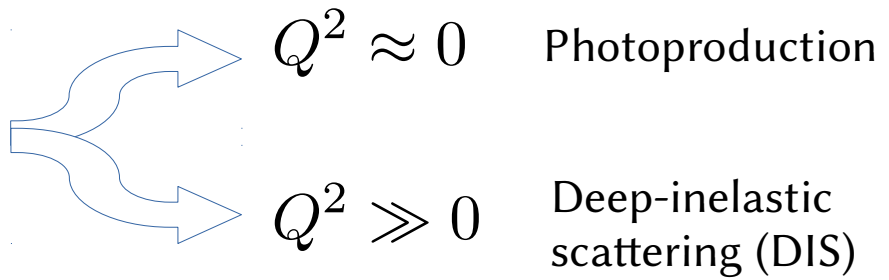


Inelasticity

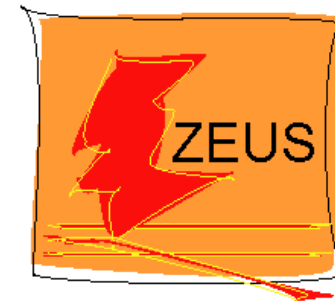
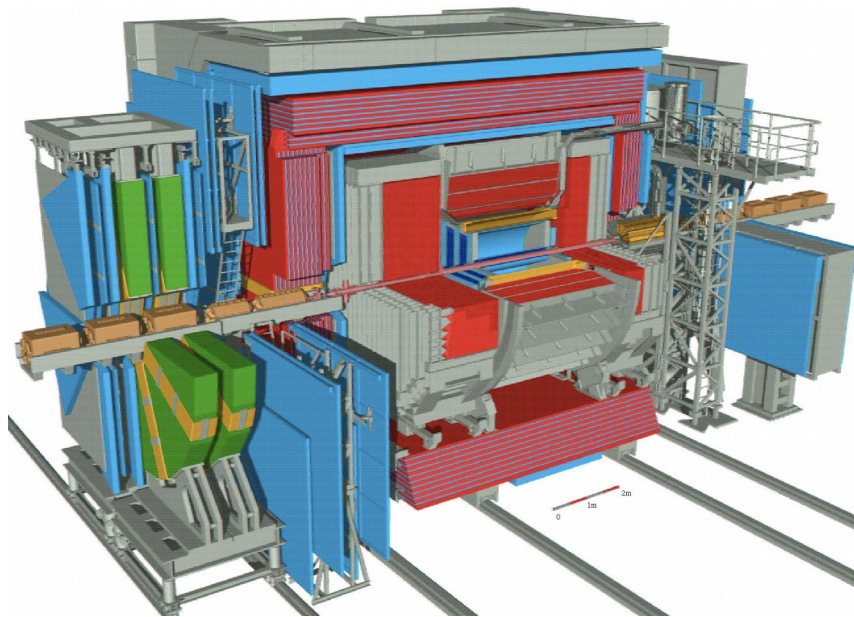
$$y = \frac{p \cdot q}{p \cdot k}$$

Photon virtuality

$$Q^2 = -(k - k')^2$$



Diffractive photoproduction of the isolated photon (ZEUS)



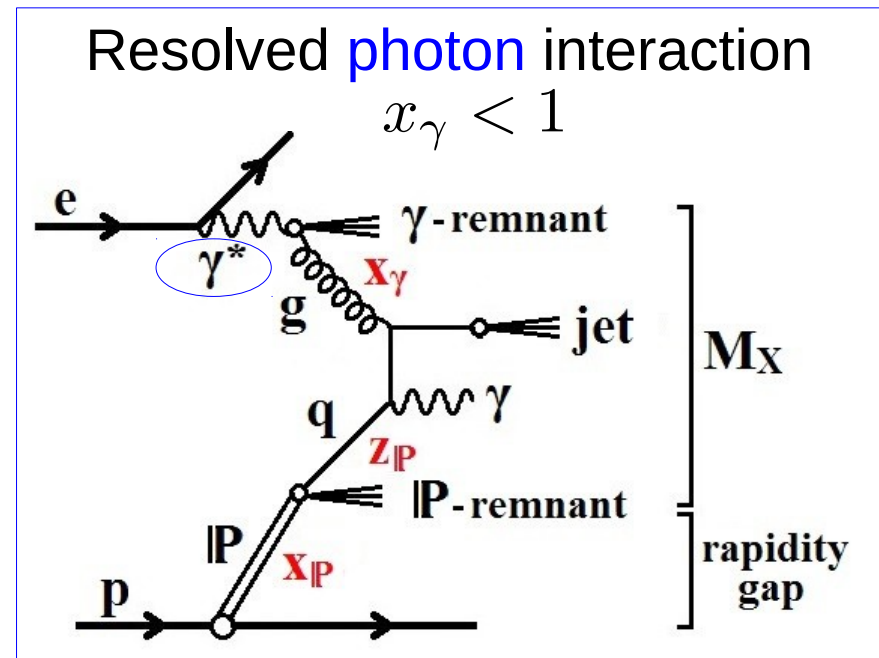
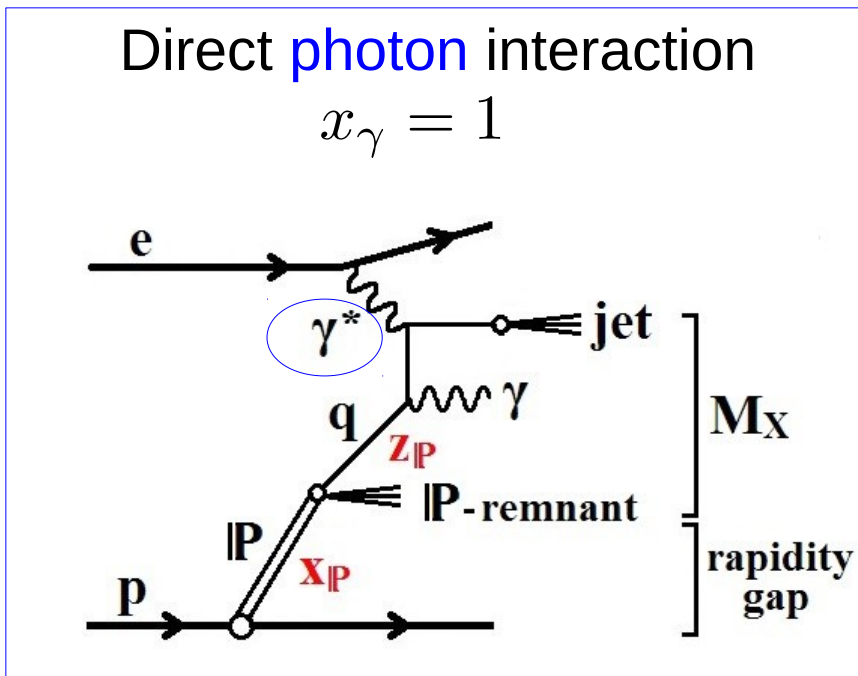
DESY-17-077 [arXiv:1705.10251]
submitted to Phys. Rev. D

Diffractive photoproduction of isolated photon

- $Q^2 \approx 0 \rightarrow$ photon may dissociate into low mass hadronic system (structure of such resolved photon described by γ PDF)
- $Q^2 \approx 0 \rightarrow \theta_e \approx 180^\circ$ (electron leaves detector undetected)

Photon momentum fraction entering the hard subprocess:

$$x_\gamma = \frac{\sum_{\gamma+\text{jet}}(E - p_z)}{\sum_{\text{EFO}}(E - p_z)}$$

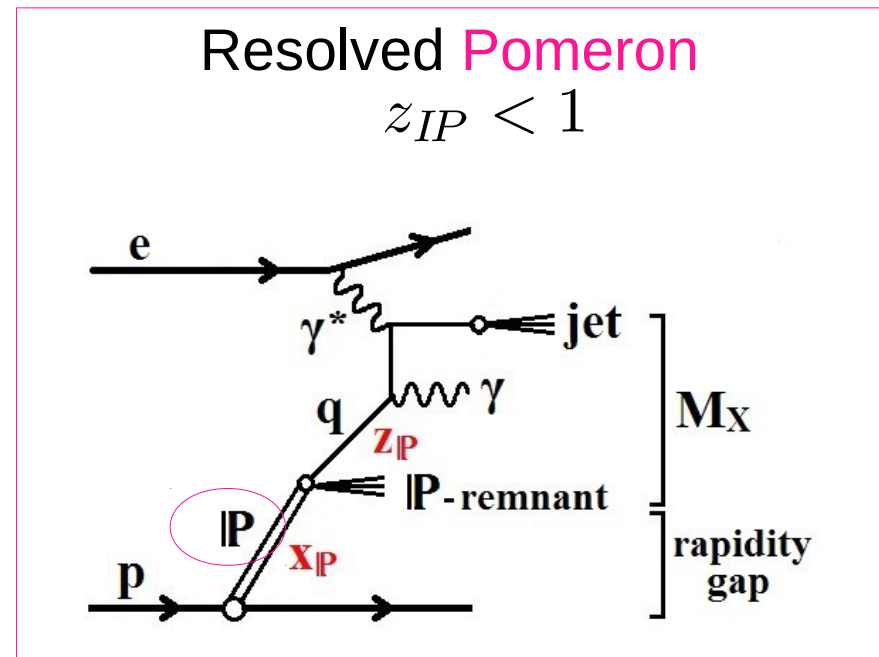
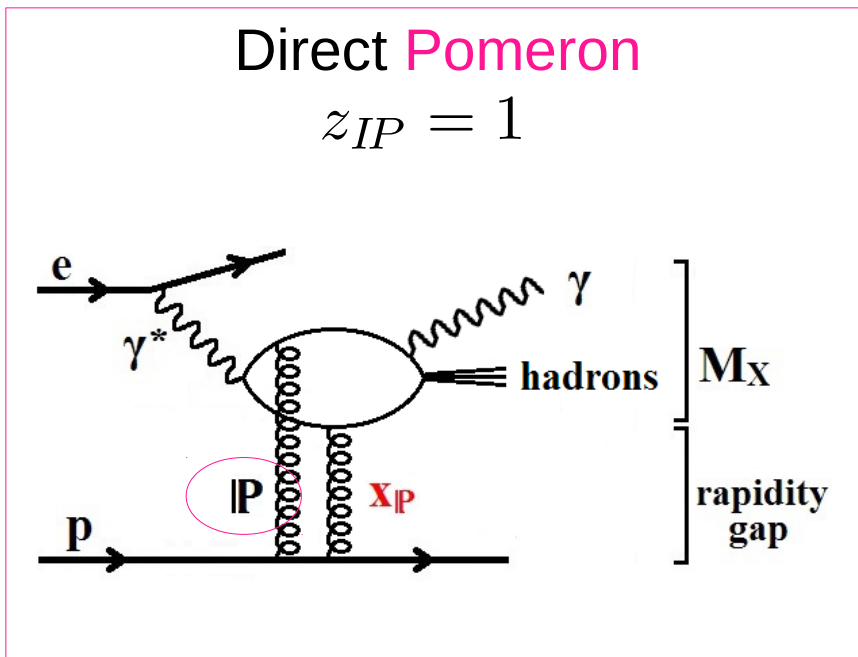


Diffractive photoproduction of isolated photon

- Diffraction \rightarrow beam proton stays intact and leaves detector undetected
- Standardly described by exchange of an hadronic object with vacuum quantum numbers (**pomeron**)

Pomeron momentum fraction entering the hard subprocess:

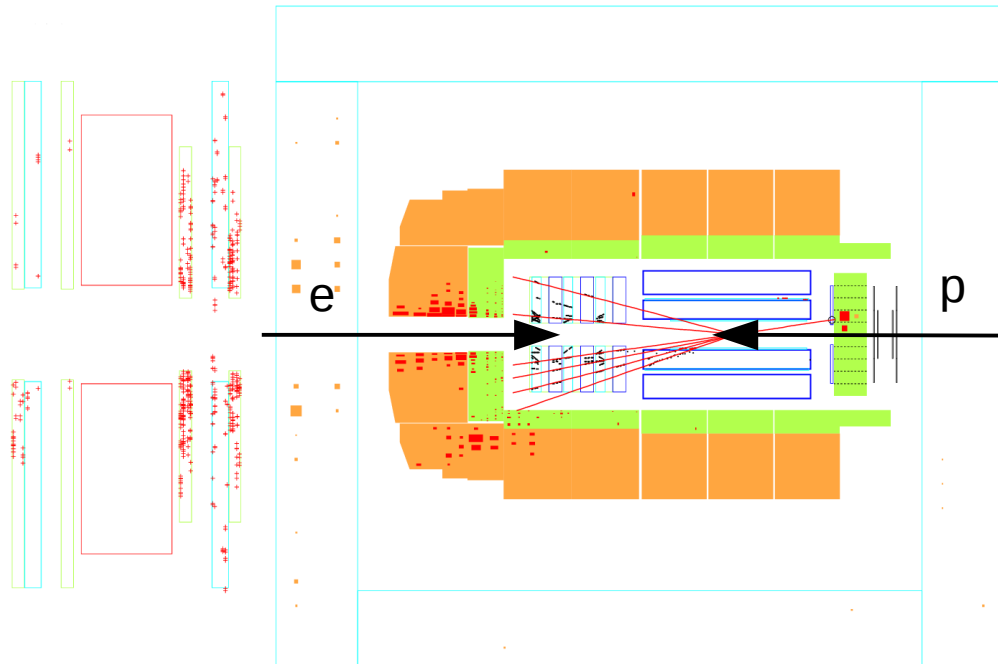
$$z_{IP} = \frac{\sum_{\gamma+\text{jet}} (E + p_z)}{\sum_{\text{EFO}} (E + p_z)}$$



Large Rapidity Gap

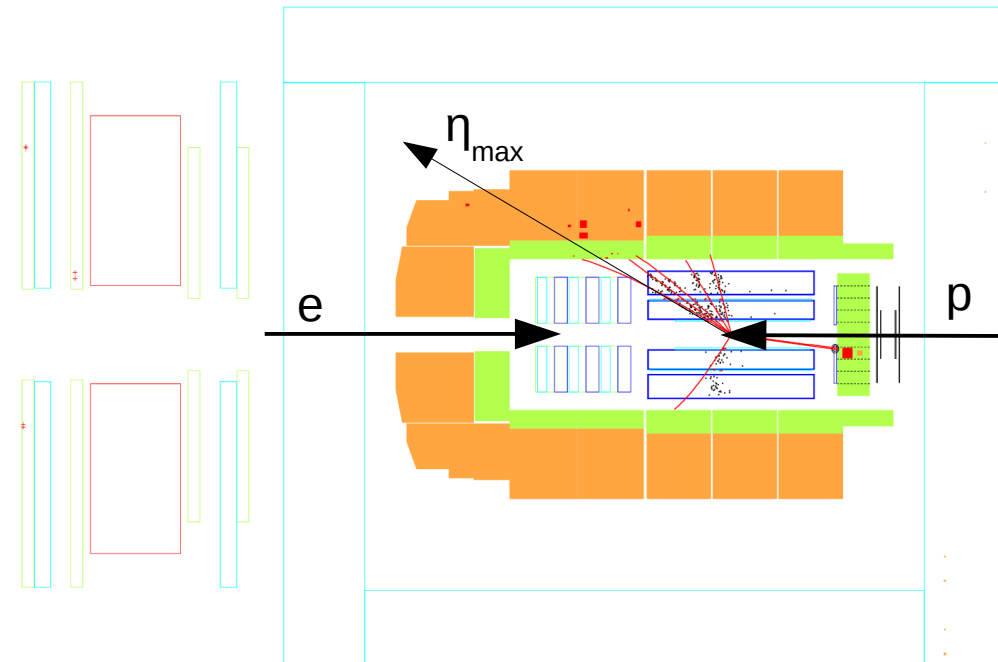
Here H1 detector

Hadronic activity in forward part of detector



Non-diffractive event

Without hadronic activity in forward part of detector



Diffractive event

Theoretical predictions

Diffractive predictions (Resolved pomeron)

- Resolved pomeron model (Ingelman Schlein)
- Implemented in the MC generator **RAPGAP** (LO matrix element + LL parton shower + Lund string fragmentation)
- Contains direct and resolved photon processes
- The partonic structure of the resolved pomeron described by **H1 2006 DPDF Fit B** (from fits of inclusive diffractive DIS)
- The partonic structure of the resolved photon described by **SASGAM-2D γ PDF**

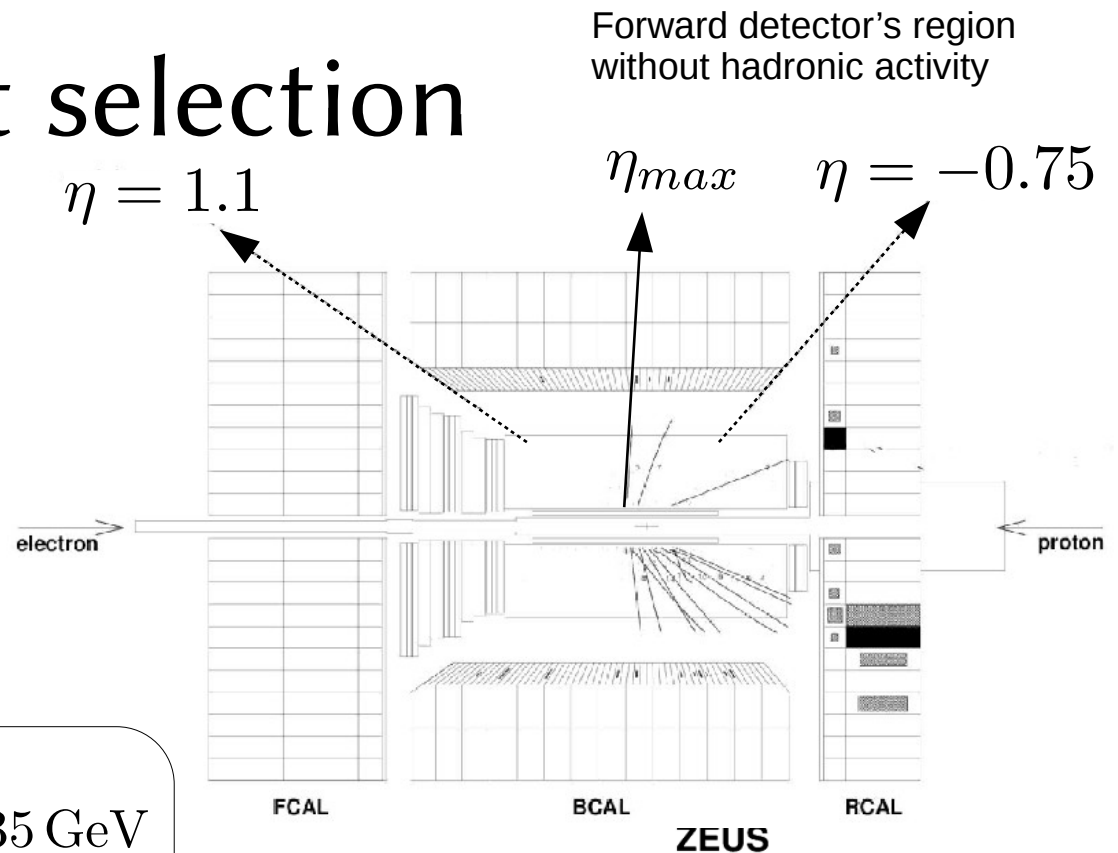
- Non-diffractive background simulated by Pythia 6

No model for the possible direct pomeron interaction available

Event selection

$$e + p \rightarrow e + \gamma + X + p(Y)$$

- Veto on scattered electron
- Diffractive events dominate for small pomeron momentum fraction wrt proton x_{IP} + **large rapidity gap**



Photon

$$5 < E_T^\gamma < 15 \text{ GeV}$$

$$-0.7 < \eta^\gamma < 0.9$$

Jet

$$4 < E_T^{jet} < 35 \text{ GeV}$$

$$-1.5 < \eta^{jet} < 1.8$$

Photoproduction

$$Q^2 < 1 \text{ GeV}^2$$

$$0.2 < y < 0.7$$

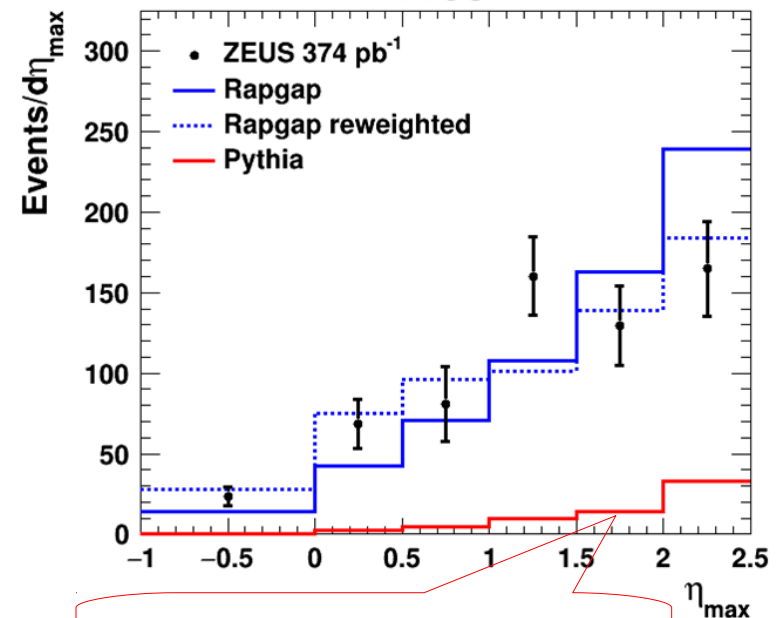
$$y = \frac{\sum_{\text{EFO}} (E - p_z)}{2E_e}$$

Diffraction

$$\eta_{max}^{E>0.4} < 2.5$$

$$x_{IP} < 0.03$$

$$x_{IP} = \frac{\sum_{\text{EFO}} (E + p_z)}{2E_p}$$



Extraction of prompt photons signal

- Template fit to obtain the signal and background contribution

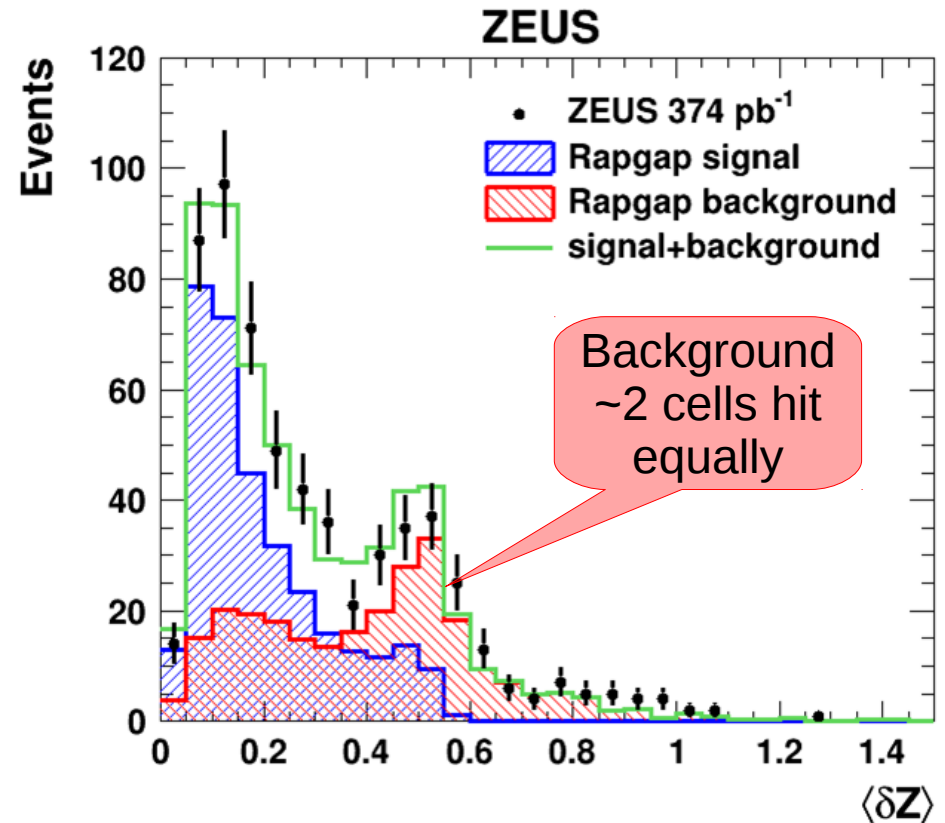
- Background mainly from

$$\pi^0(\eta) \rightarrow \gamma\gamma$$

- Width of the photon candidate cluster in the beam direction in units of cell width δ_{cell}

$$\langle \delta Z \rangle = \frac{\sum_i E_i |Z_i - Z_{cluster}|}{\delta_{cell} \sum_i E_i}$$

- **90%** of photon candidate energy required to be measured in EM calorimeter



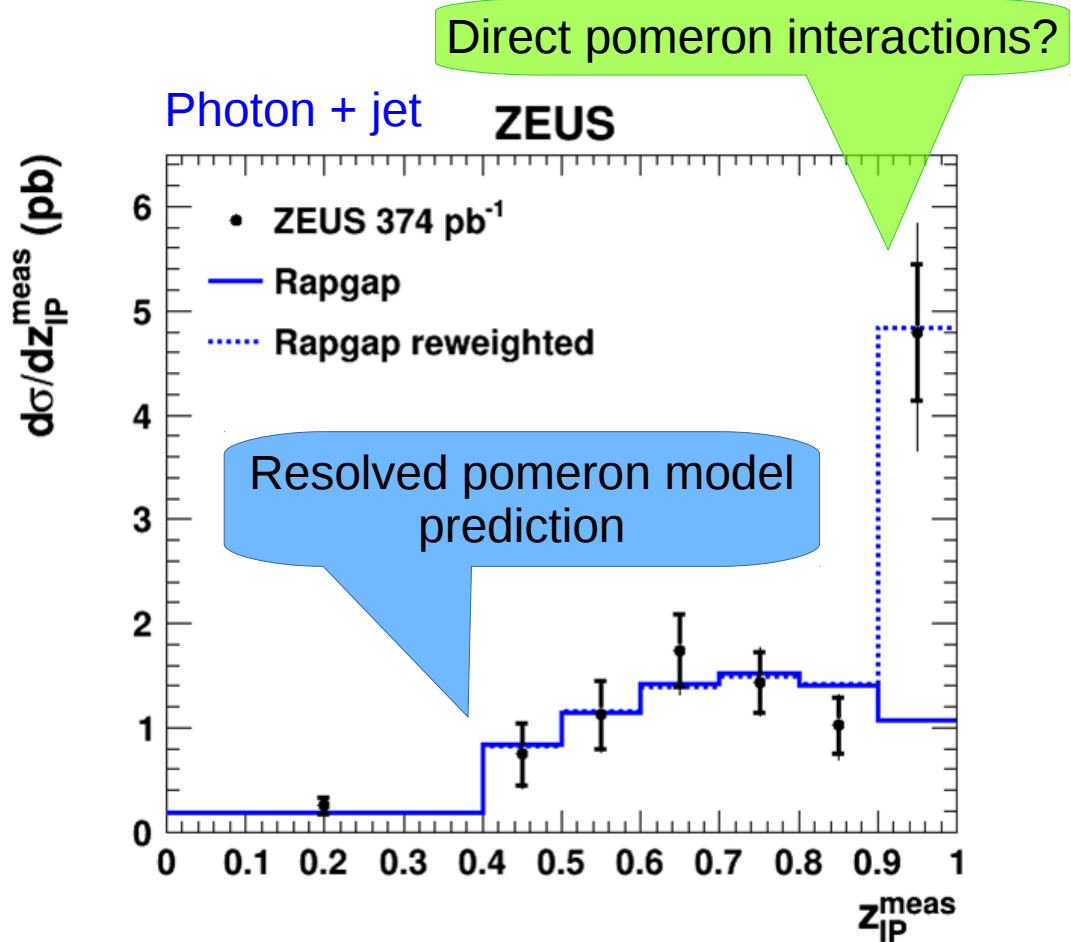
	Gamma events	Gamma+jet events
HERA I (82 pb ⁻¹)	91	76
HERA II (374 pb ⁻¹)	366	311

Direct pomeron exchange?

- **The $z_{IP} < 0.9$ region** well described by MC both in shape and normalization

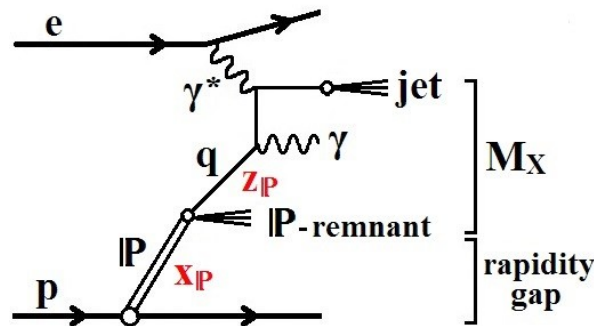
$$\sigma_{\text{data}}^{z_{IP} < 0.9} = 0.57 \pm 0.13 \text{ pb}$$

$$\sigma_{\text{MC}}^{z_{IP} < 0.9} = 0.68 \text{ pb}$$
- **The $z_{IP} > 0.9$ region** overshoot in data

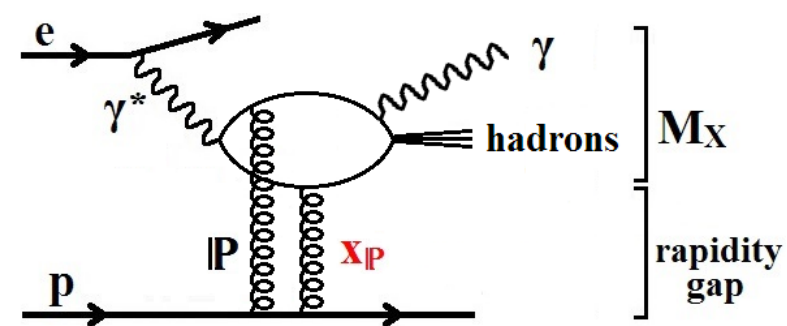


Rapgap reweighted:
MC reweighted separately for $z_{IP} < 0.9$ and $z_{IP} > 0.9$ to data

Resolved pom.



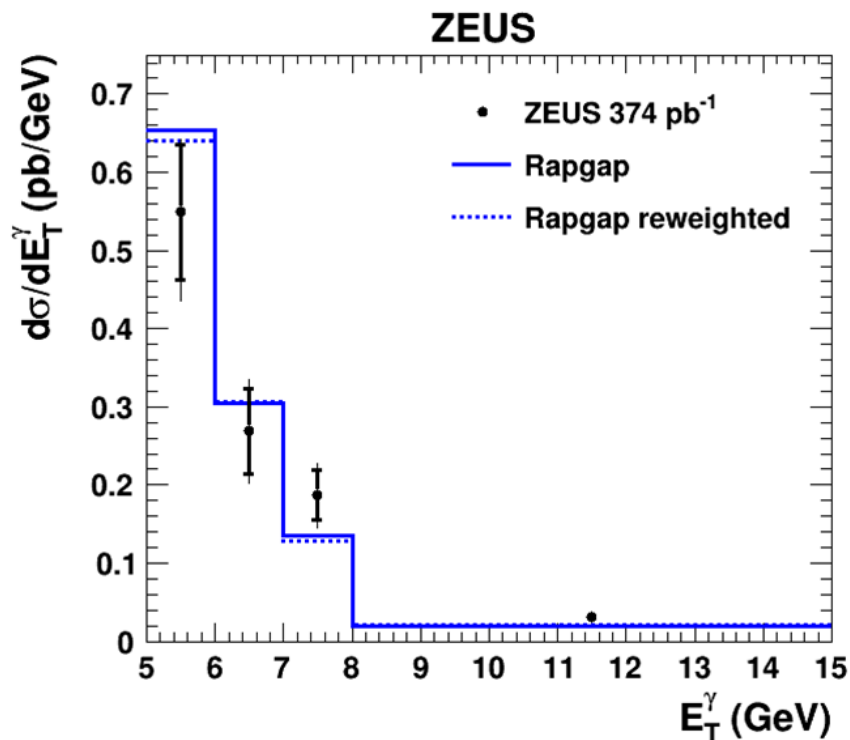
Direct pom.



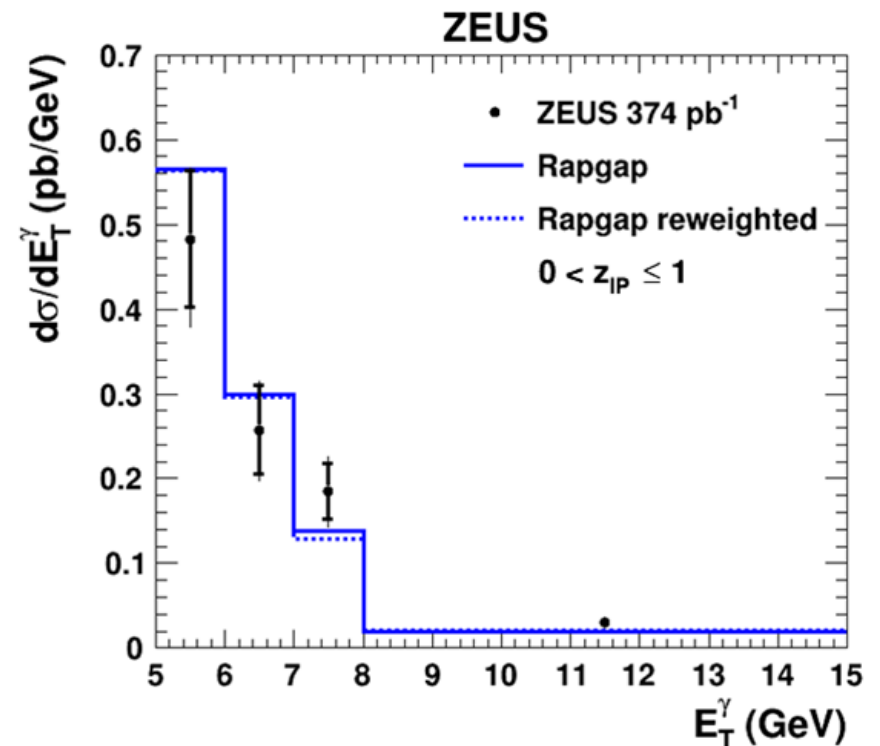
Prompt photon transverse momentum

- Shape of the gamma transverse momentum well described by MC prediction (MC always **normalized to data**)
- 85% of events with prompt photon contain jet as well

Prompt photons



Prompt photons with jet



Spectrum of x_{IP}

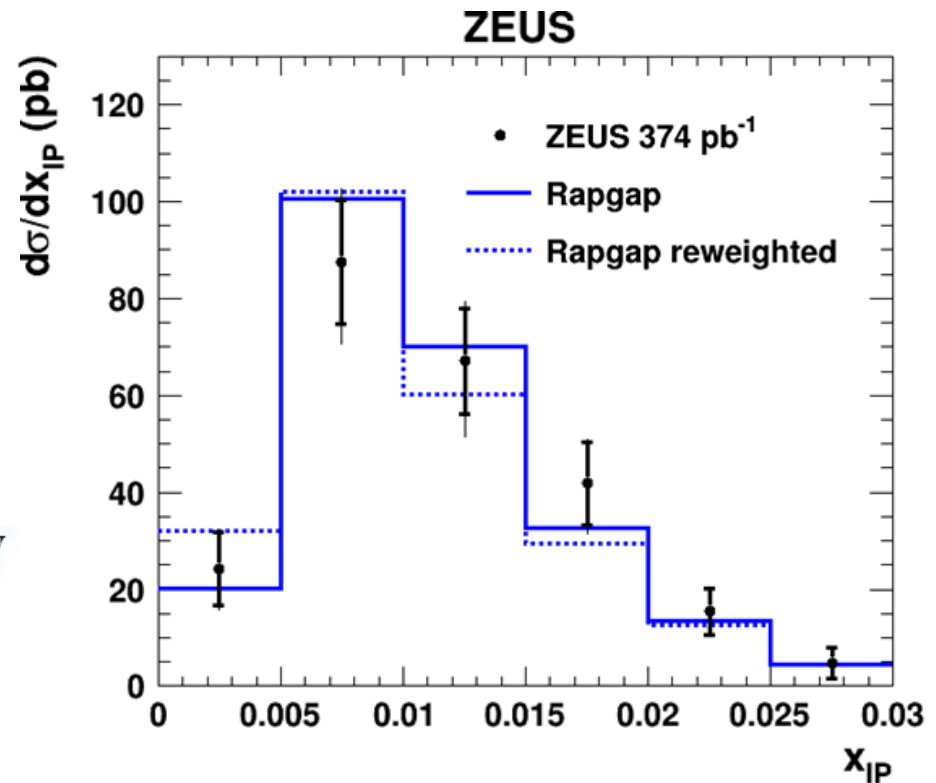
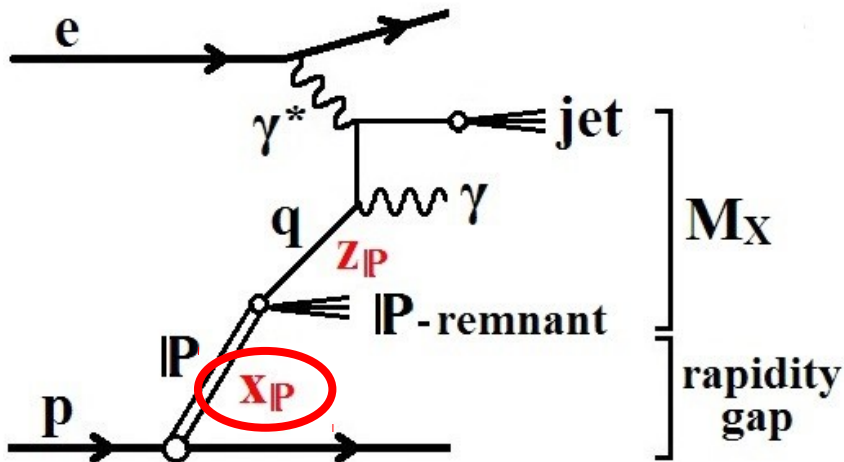
- The relative energy loss of the leading proton with respect to the incoming beam proton ($\sim 1\%$)
- As leading proton directly not measured, reconstructed from EFO

$$x_{IP} = 1 - \frac{E'_p}{E_p}$$

↕

$$x_{IP} = \frac{\sum_{\text{EFO}} (E + p_z)}{2E_p}$$

Prompt photons



Direct vs Resolved pomeron

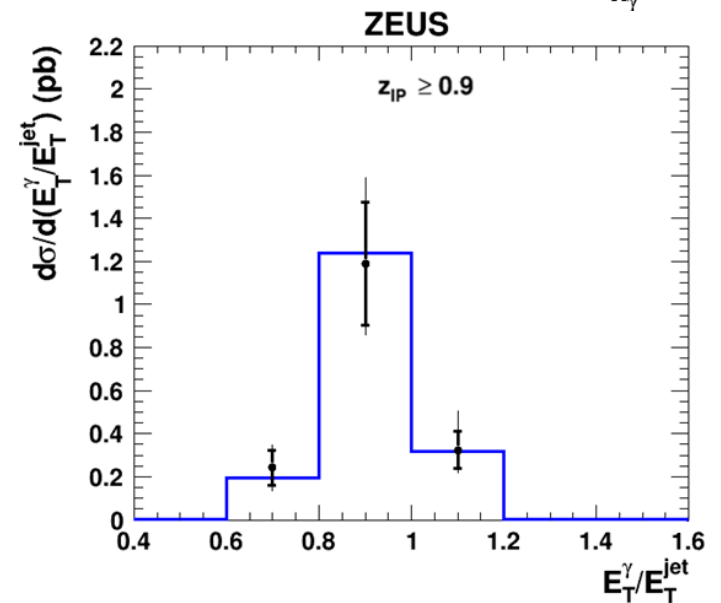
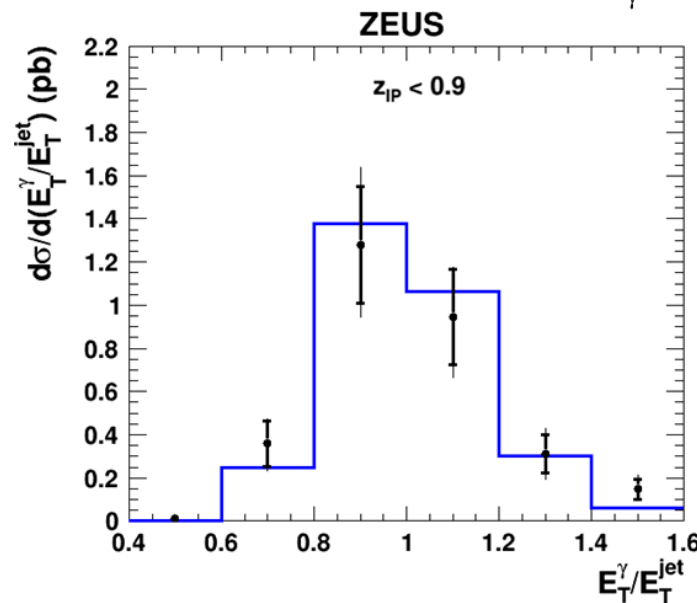
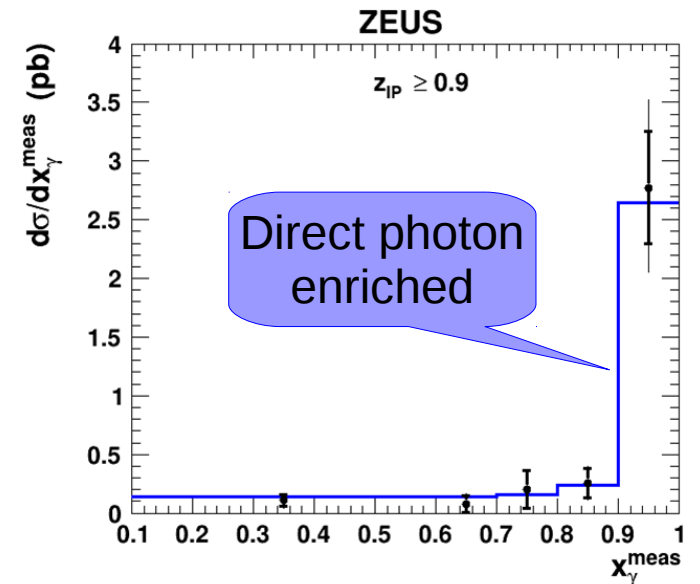
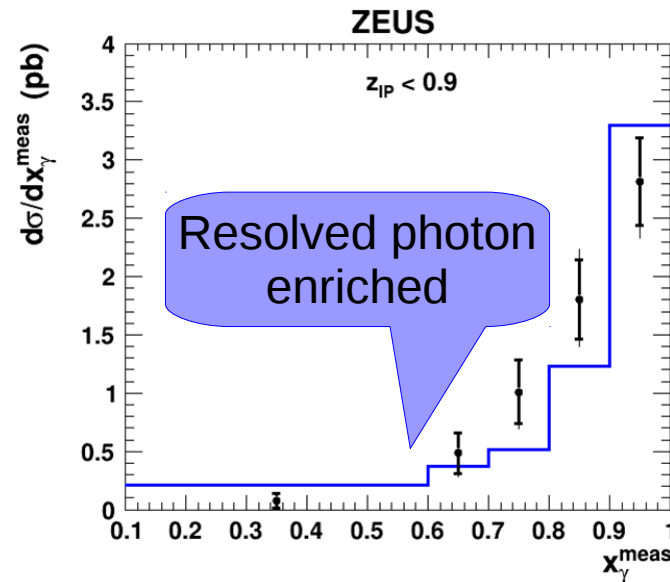
Photon + jet

Photon momentum fraction entering hard process

Gamma/jet p_T balance

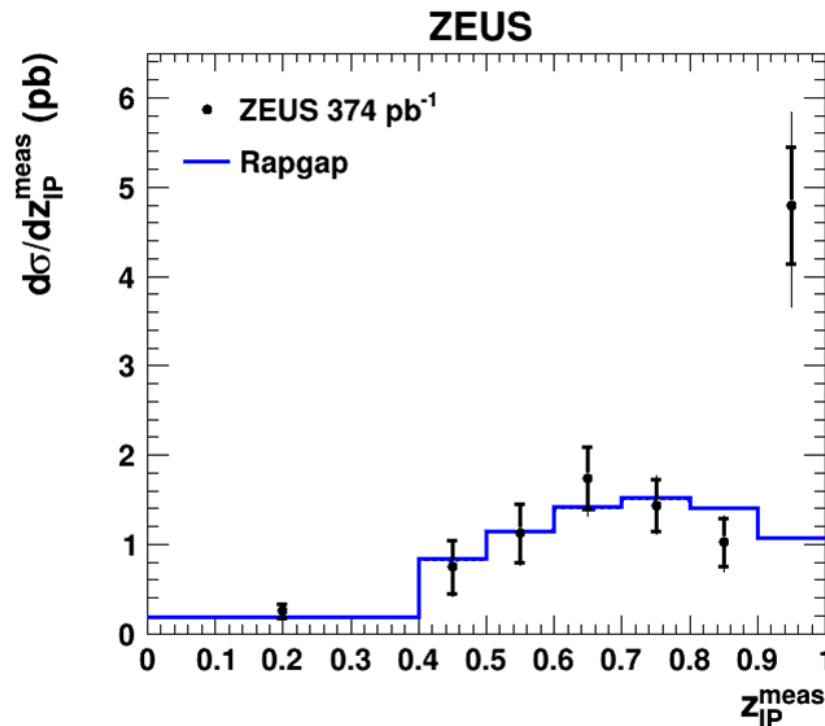
Resolved pomeron enriched ($z_{IP} < 0.9$)

Direct pomeron enriched ($z_{IP} \geq 0.9$)

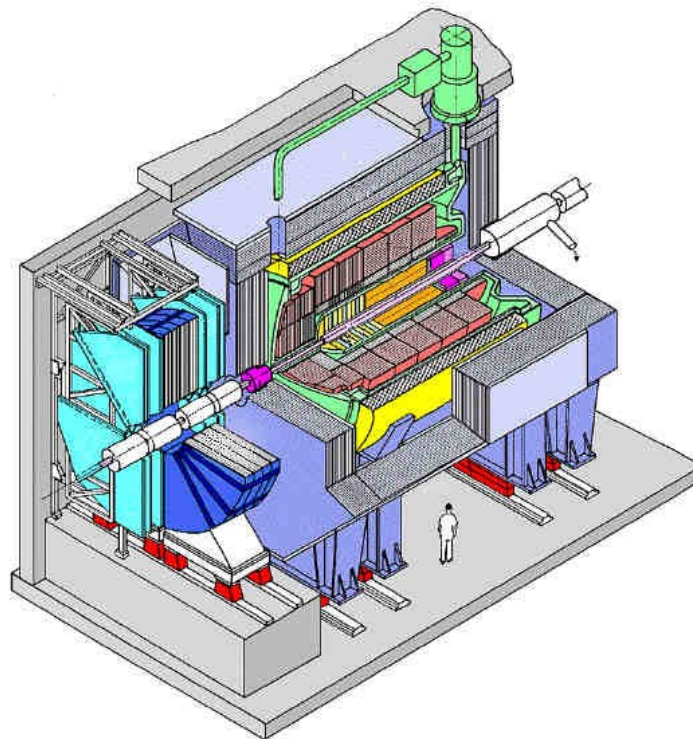


Conclusion 1

- The first measurement of diffractively produced prompt photon
- The peak in z_{IP} distribution suggests direct pomeron interactions (occurring mostly in direct photon interactions)
- Other distributions well described by MC prediction



Extraction of α_S at NNLO from jet cross sections in DIS (H1)



H1prelim-17-031

[http://www-h1.desy.de/publications/H1preliminary.short_list.html]

DESY-16-200 Eur.Phys.J.C77 (2017) 4, 215 [arxiv:1611.03421]

NNLO α_S fit of H1 jets data in DIS

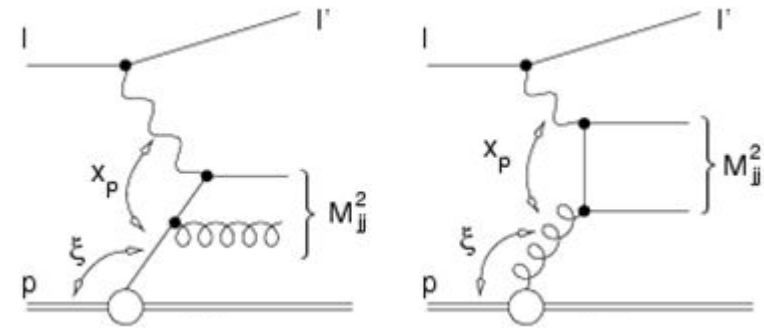
Why α_S ?

- Among the least known SM parameters

$$G_F = 1.1663787(6) \times 10^{-5} \text{ GeV}^{-2}$$

$$\alpha_S = 0.1181(11) \text{ [PDG16]}$$

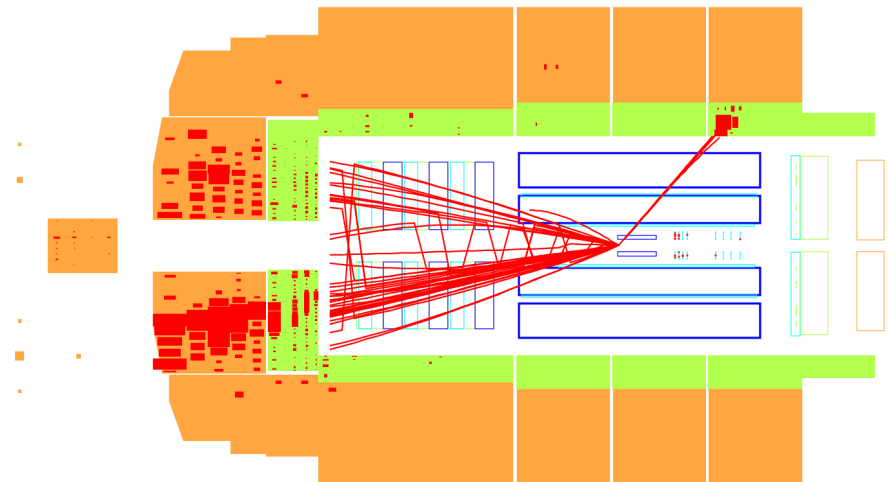
- Great importance for LHC physics



Dijet DIS production at H1

Why now?

- NNLO revolution in the last years
- NNLO predictions now available for both pp and ep dijets
- LHC has not fitted their **jet** data with NNLO yet

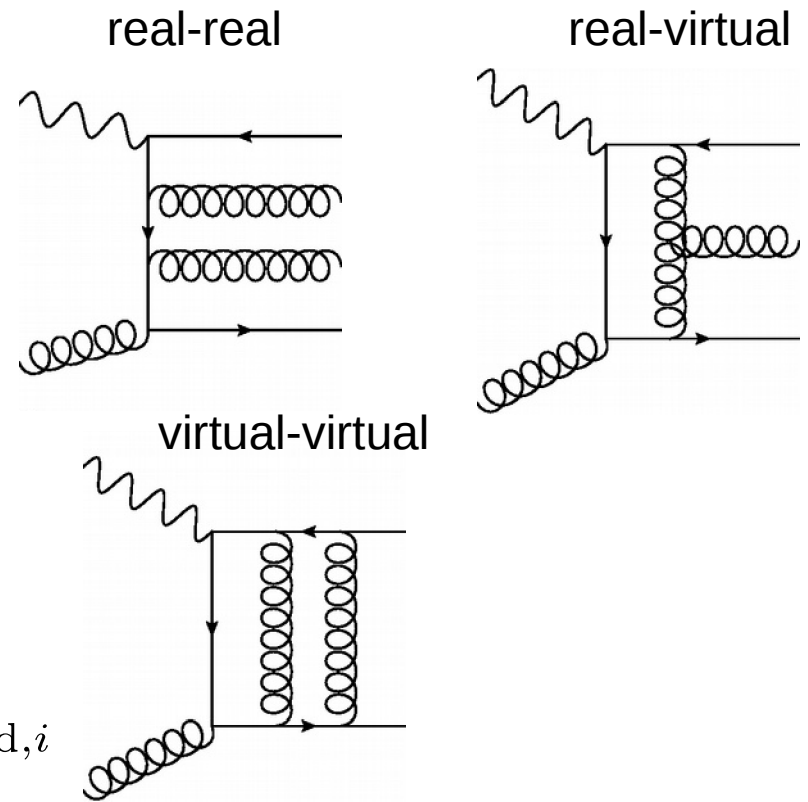


First NNLO α_S fit of the jet ep data

NNLO calculations

✓ Real-real + real-virtual crosschecked with NLOJET++ & SHERPA

- New NNLO predictions for ep dijets based on **antenna subtraction**
J. Currie, T. Gehrmann, A. Huss and J. Niehues, JHEP 07 (2017) 018, [1703.05977]
- **Matrix element** tables precalculated by **NNLOJET** program (~1M CPU hours)
- Then convoluted with PDFs and α_S using **fastNLO** (<1s)



$$\sigma_i = \sum_{k=g,q,\bar{q}} \int dx f_k(x, \mu_F) \hat{\sigma}_{i,k}(x, \mu_R, \mu_F) \cdot C_{\text{had},i}$$

fastNLO

$$\hat{\sigma}_{i,k}(x, \mu_R, \mu_F) =$$

$$\sum_n \alpha_S^n(\mu_R) \hat{\sigma}_{i,k}^{(n)}(x, \mu_R, \mu_F)$$



A bit of history

- **1973** Asymptotic freedom of QCD
- **1993** NLO studies of DIS jets
- **2016** NNLO corrections for DIS jets

H1 Data – Dijets

Double-diff.

- Q^2 and $\langle p_T \rangle_2$
- Mean dijet p_T

$$\langle p_T \rangle_2 = \frac{p_T^{\text{jet1}} + p_T^{\text{jet2}}}{2}$$

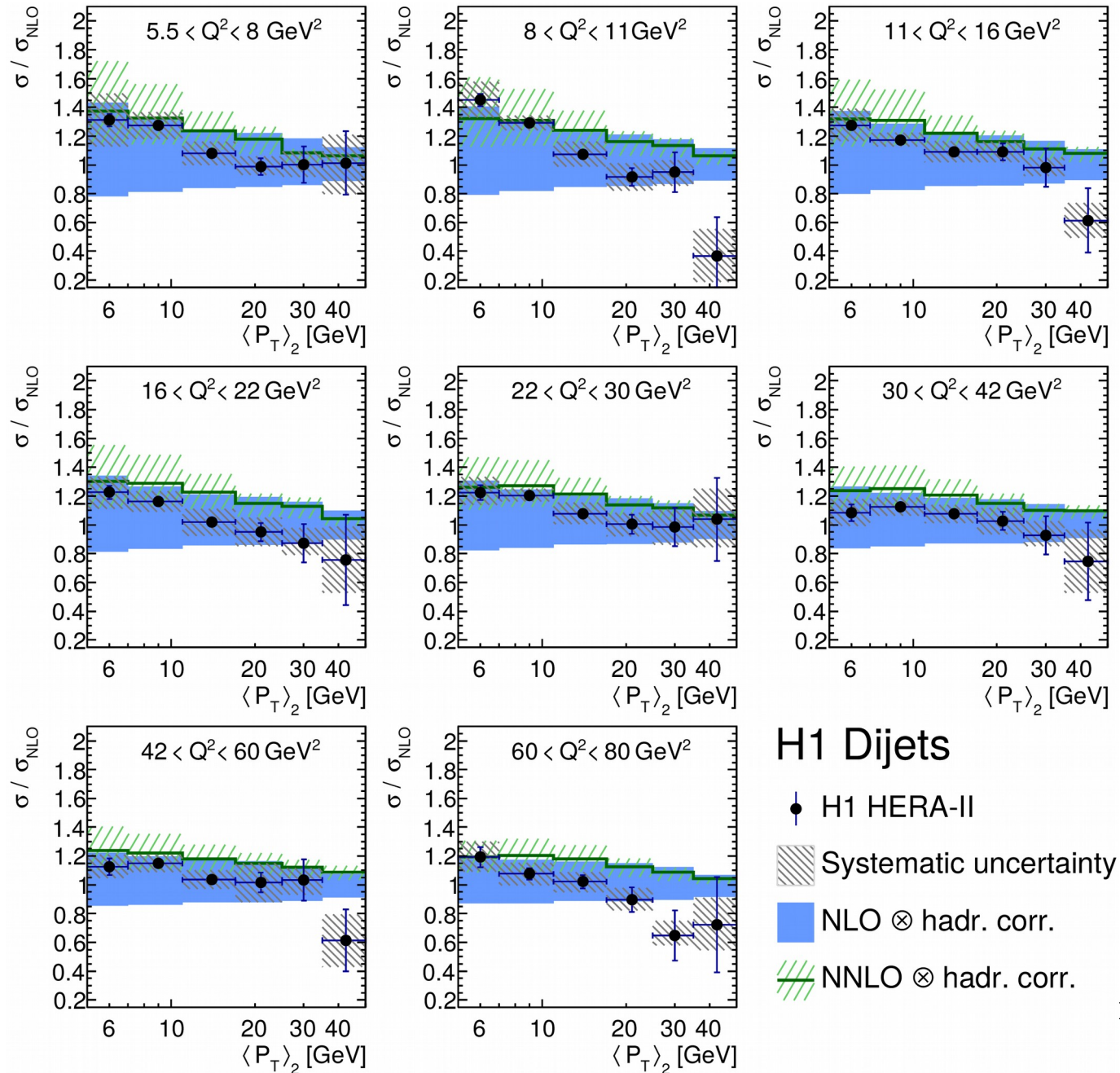
jets found in $\gamma^* p$
with k_T algo ($R=1$)

NLO predictions

- NNPDF 3.0 NLO
- Larger scale unc.
- $\text{Chi2}/\text{ndf} = 1.4$

NNLO predictions

- NNPDF 3.0 NNLO
- Smaller scale unc.
- $\text{Chi2}/\text{ndf} = 0.6$



H1 Dijets

- H1 HERA-II
- Systematic uncertainty
- NLO \otimes hadr. corr.
- NNLO \otimes hadr. corr.

H1 Data – Inclusive jets

Double-diff.

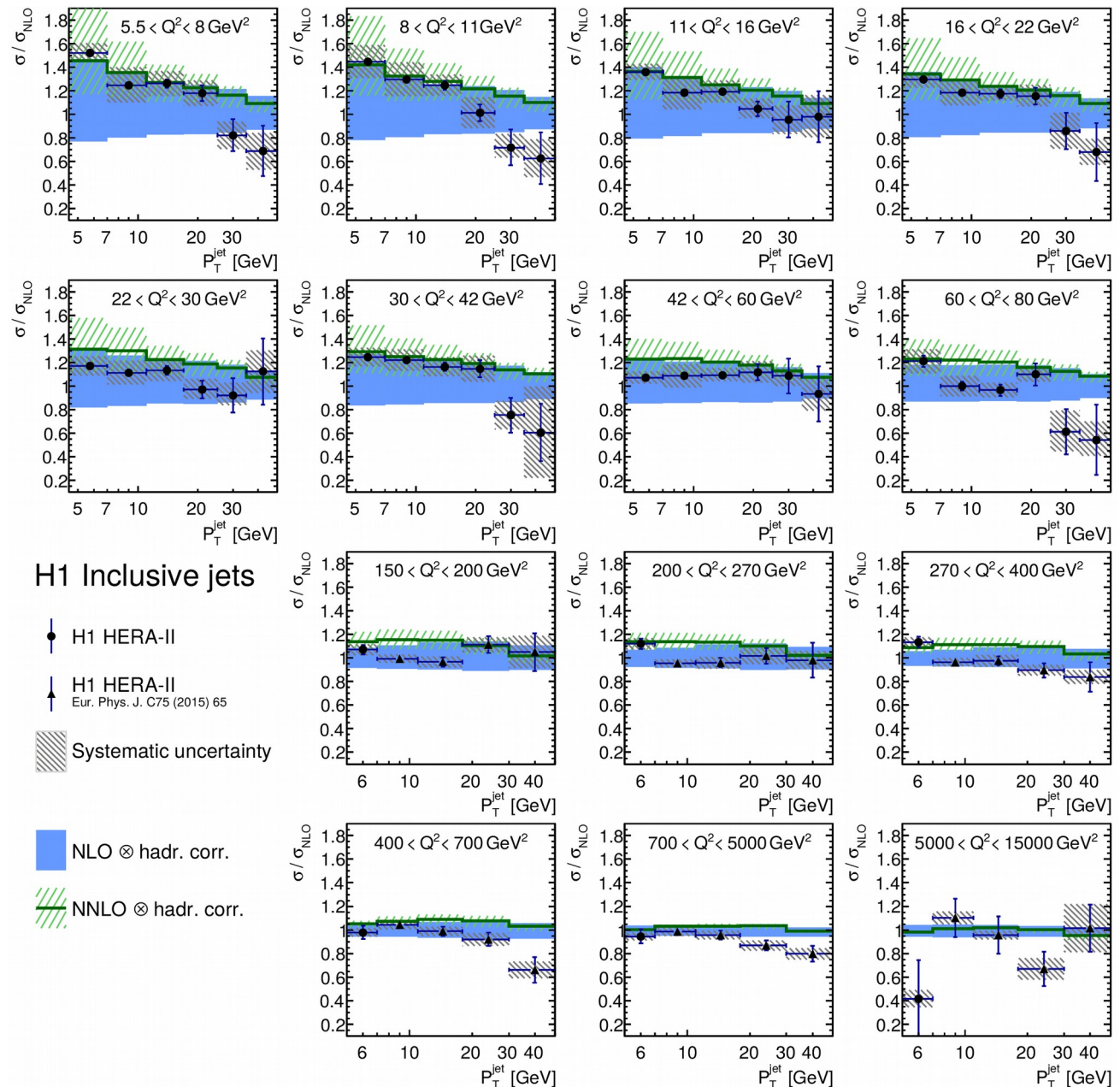
- Q^2 and p_T^{jet}
- Mean dijet p_T
- $0.2 < y < 0.6$
- $-1 < \eta_{\text{lab}}^{\text{jet}} < 2.5$
- jets found in $\gamma^* p$ with k_T algo ($R=1$)

NLO predictions

- *NNPDF 3.0 NLO*
- Larger scale unc.
- $\text{Chi}2/\text{ndf} = 1.7$

NNLO predictions

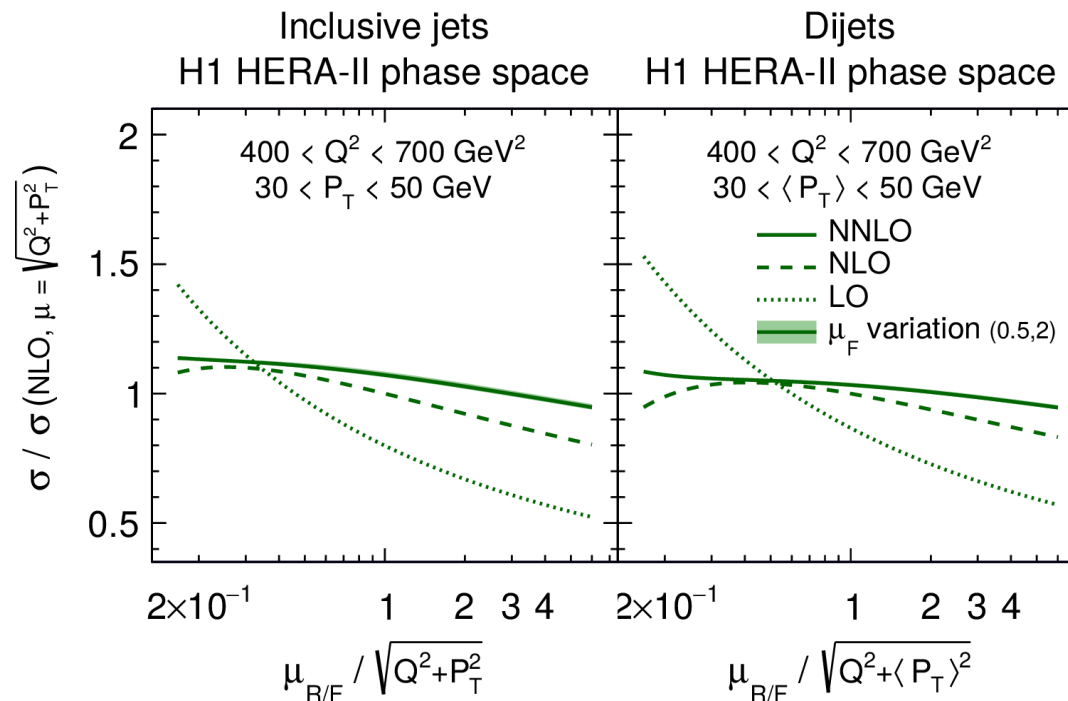
- *NNPDF 3.0 NNLO*
- Smaller scale unc.
- $\text{Chi}2/\text{ndf} = 1.3$



Scale dependence

- The NNLO predictions depend **less** on the renormalization scale (=have smaller theor. unc.)
- To estimate the uncertainty the scale varied up and down by the factor of 2
- As a scale we use $\mu_R = \mu_F = \sqrt{Q^2 + p_T^2}$

Others functional forms also tested



Functional form of the scale

- 7 possible function studied
- NNLO α_S is usually smaller than the NLO one
- The NNLO χ^2 is usually better
- NNLO scale unc. is smaller

$$\mu^2 = Q^2$$

$$\mu^2 = p_T^2$$

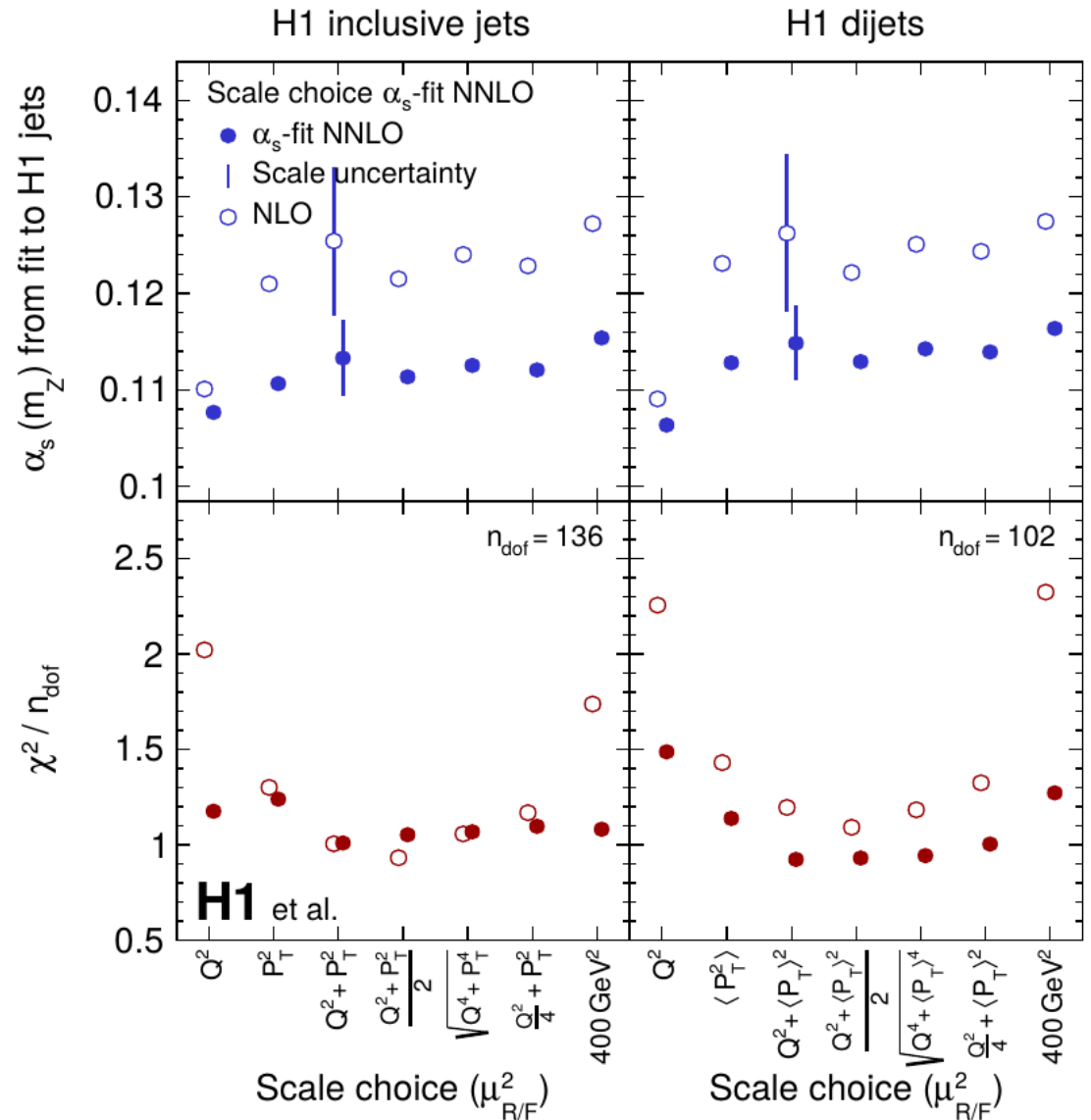
$$\mu^2 = Q^2 + p_T^2$$

$$\mu^2 = \frac{Q^2 + p_T^2}{2}$$

$$\mu^2 = \sqrt{Q^4 + p_T^4}$$

$$\mu^2 = Q^2/4 + p_T^2$$

$$\mu^2 = 400 \text{ GeV}^2$$



All data above m_b threshold used

α_S in PDF and α_S in ME

- Alpha strong affecting both, PDFs and matrix element
- Both effects considered, α_S in ME more prominent

$$\sigma_i = \sum_{k=g,q,\bar{q}} \int dx f_k(x, \mu_F) \hat{\sigma}_{i,k}(x, \mu_R, \mu_F) \cdot C_{\text{had},i}$$

α_S dep. α_S dep.

$$\hat{\sigma}_{i,k}(x, \mu_R, \mu_F) = \sum_n \alpha_S^n(\mu_R) \hat{\sigma}_{i,k}^{(n)}(x, \mu_R, \mu_F)$$

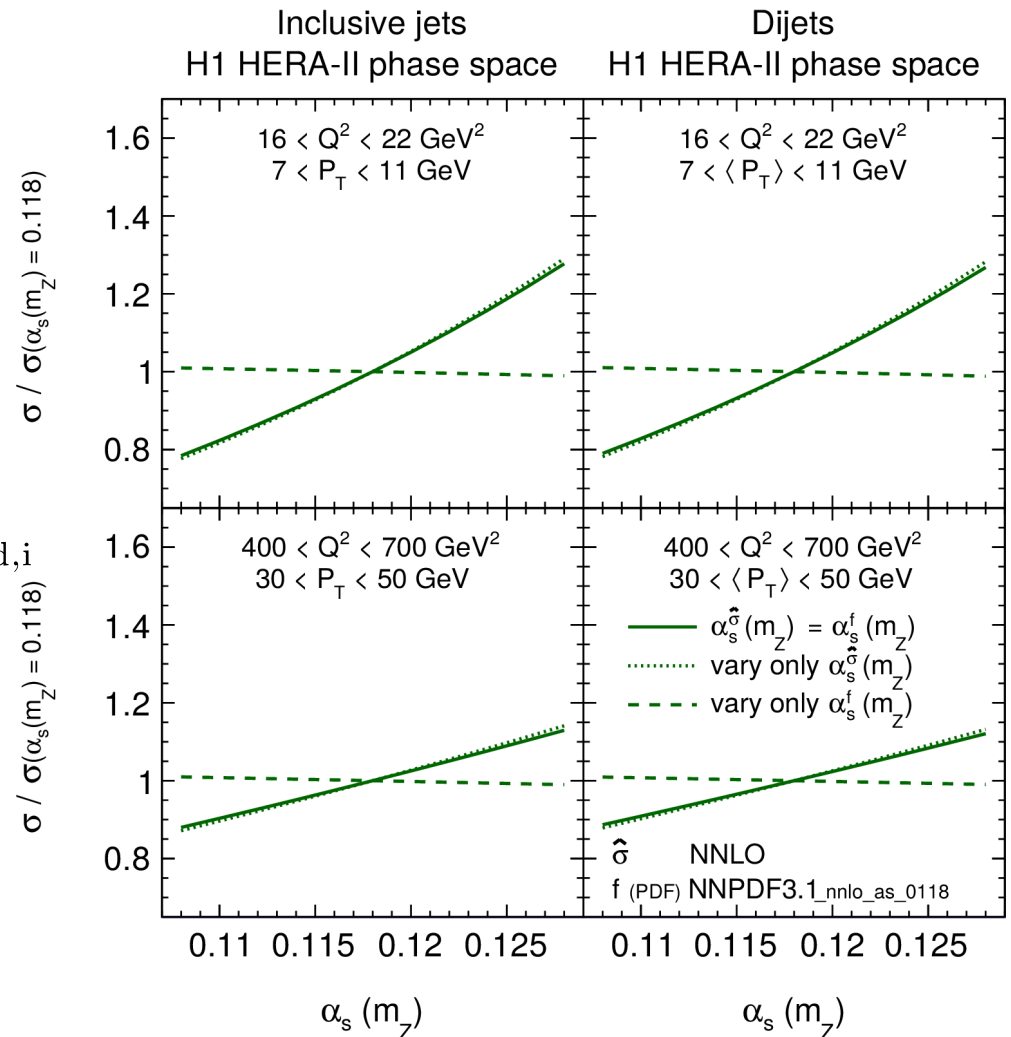
DGLAP equations

$$\mu_F^2 \frac{df}{d\mu_F^2} = P(z, \alpha_S) \otimes f(x, \mu_F^2)$$

PDFs at scale $\mu_0 = 30 \text{ GeV}$ very well constrained by lot of data
 $\rightarrow \alpha_S$ - “independent”



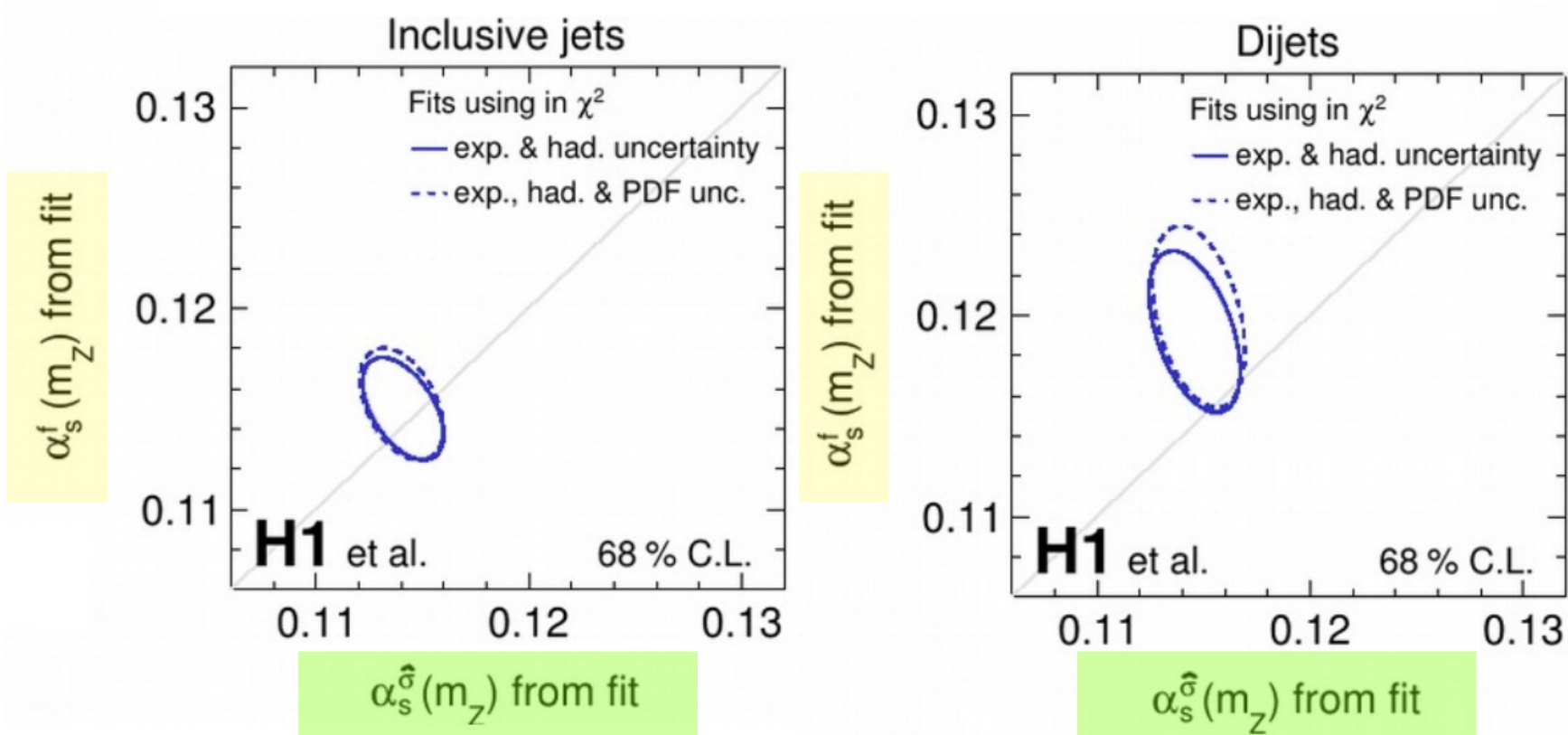
Original PDFs from scale $\mu_0 = 30 \text{ GeV}$ evolved to higher/lower scales by DGLAP with $\alpha_S = \alpha_S(\text{fit par.})$



Independent fitting of two α_S

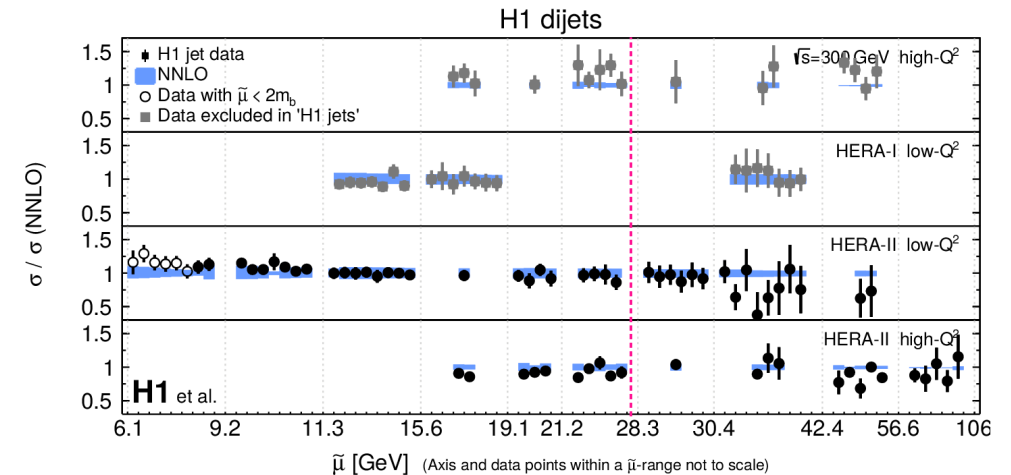
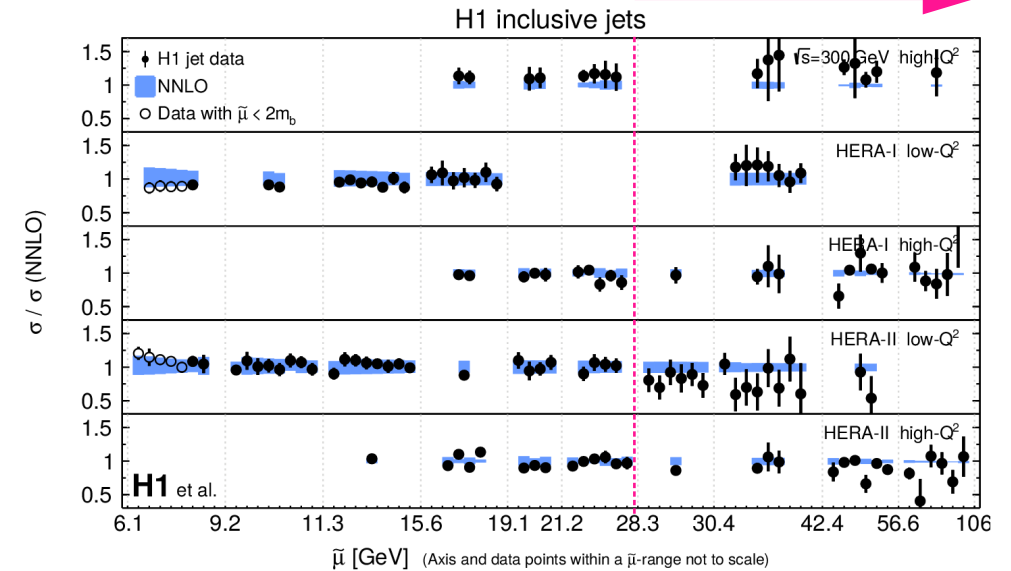
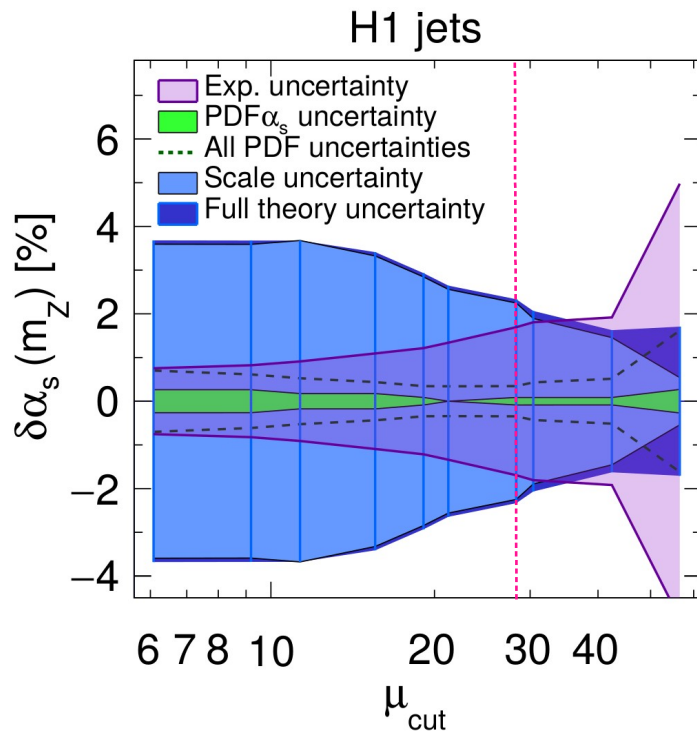
- The alpha strong from PDF and ME consistent

$$\sigma_i = f\left(\alpha_s^f(m_Z)\right) \otimes \hat{\sigma}_i\left(\alpha_s^{\hat{\sigma}}(m_Z)\right) \cdot C_{\text{had},i}$$



Which data use in the fit?

- The scale uncertainty gets higher with smaller scales
 $(\mu = \sqrt{p_T^2 + Q^2})$
- We use only data $\mu > \mu_{cut}$

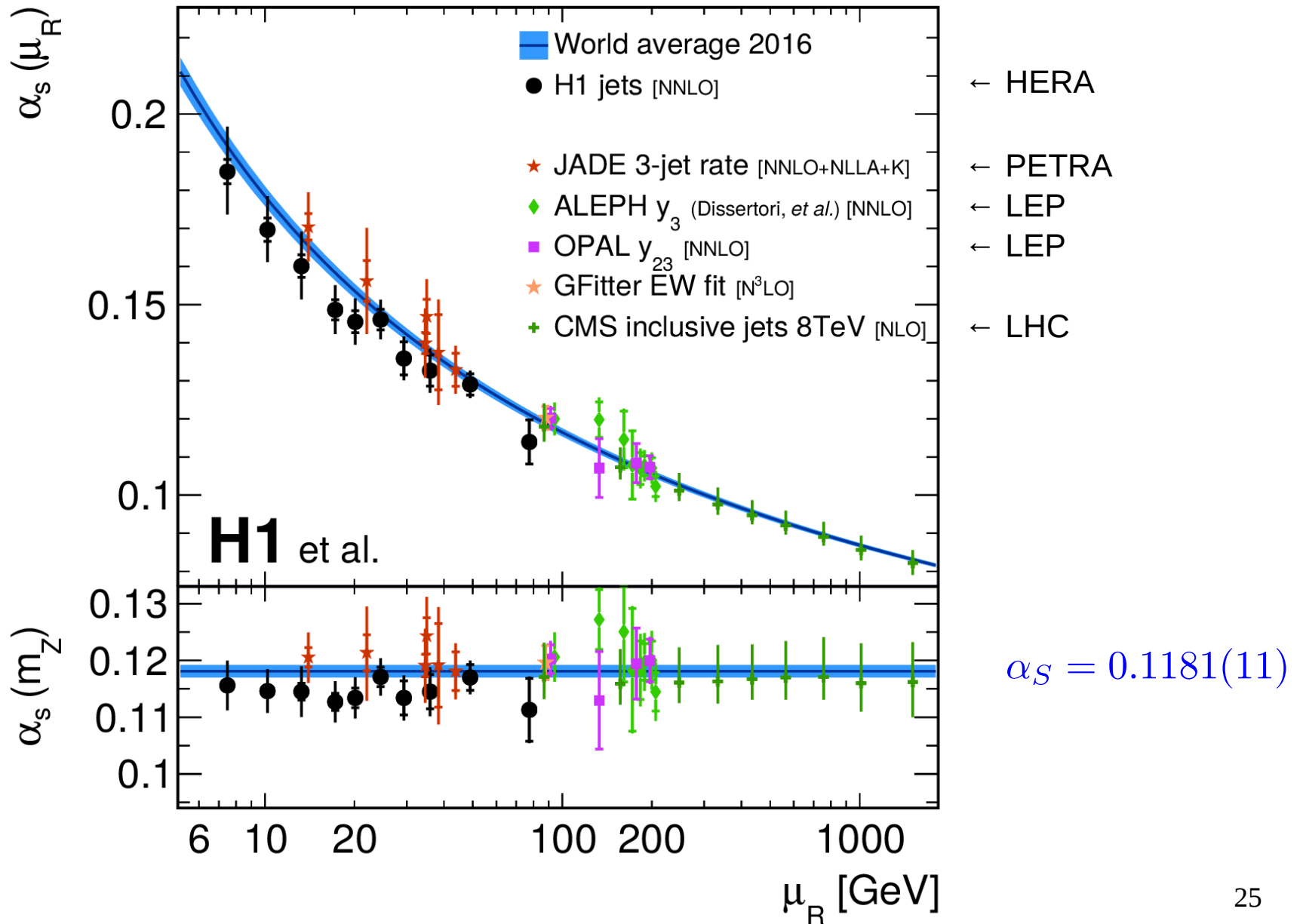


Small μ_{cut} \rightarrow high theor. unc.
 Large μ_{cut} \rightarrow high exp. unc.



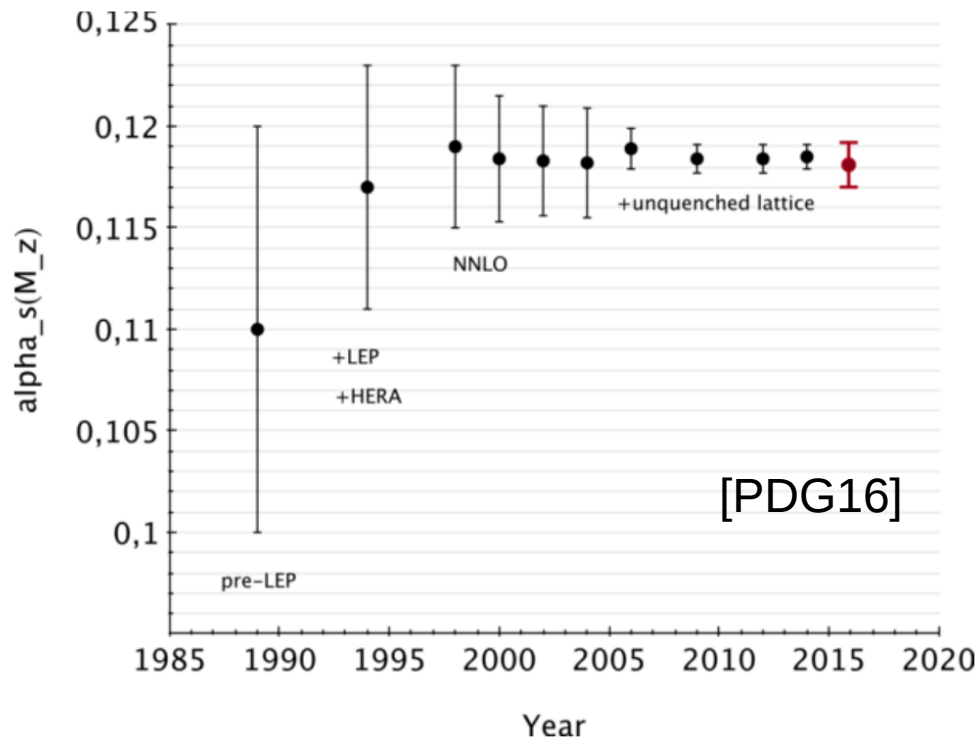
Compromise $\mu_{cut} = 28$ GeV

Running of alpha strong

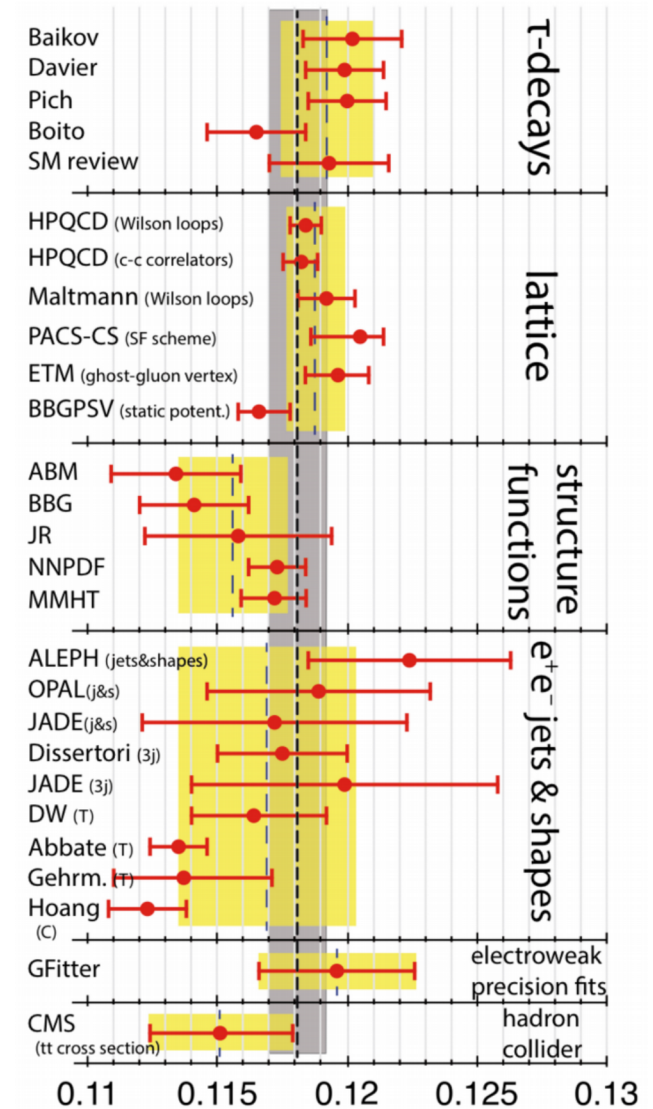


History of Alpha Strong

- The current world average value $\alpha_s(m_Z) = 0.1181(11)$
- Mostly driven by lattice and tau-decays
- From LHC the most precise estimate is from $t\bar{t}$ (NNLO)



At least NNLO fits [PDG16]

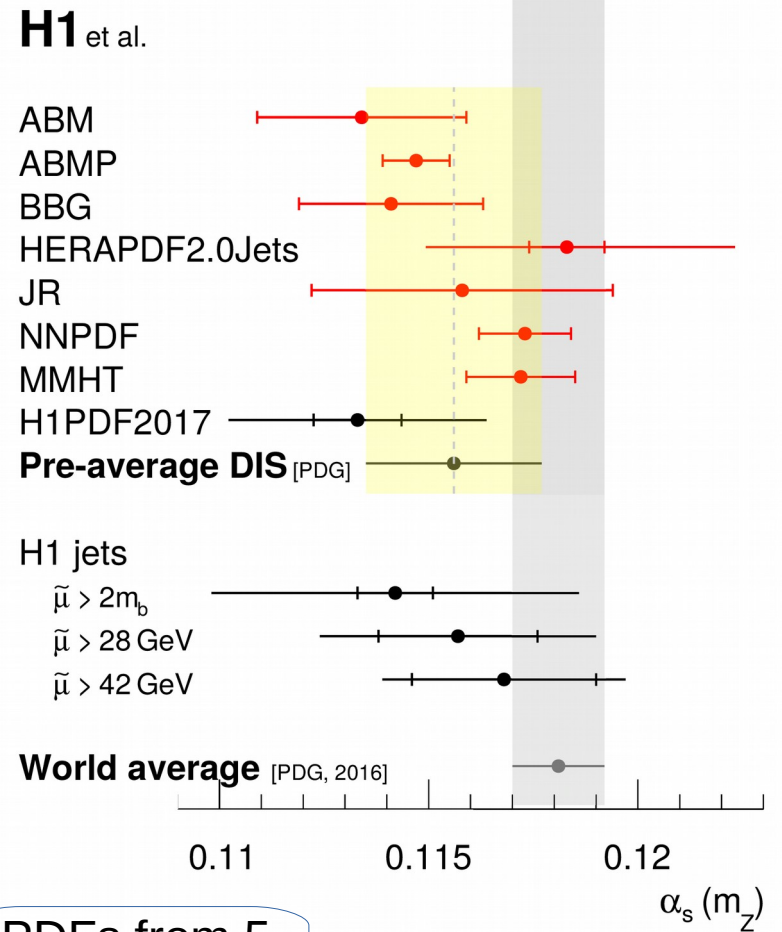


April 2016

$\alpha_s(M_Z^2)$ 26

Measured alpha strong value

- Consistent with the “world average value”
- Consistent with α_S from global PDF fits
- The NNLO reduces the scale uncertainty **by half**
- The theoretical scale uncertainty still dominant



Hadronisation

PDFs from 5 collaborations

$$\alpha_S(m_Z) = 0.1156(19)_{\text{exp}}(6)_{\text{had}}(3)_{\text{PDF}}(1)_{\text{PDF}\alpha_S}(3)_{\text{PDFSet}}(25)_{\text{scale}}$$

Data unc.

NNPDF 3.1
unc.

NNPDF 3.1
 α_S variants

Scale unc.

Conclusion 2

- The α_S from the jet DIS data estimated with NNLO precision for the first time
- The obtained value competitive with LHC and LEP measurements
- The uncertainty of H1 data even now smaller than the theoretical one \rightarrow waiting for N³LO

S. Bethke, Nucl.Part.Phys.Proc. 282-284 (2017) 149-152

Subclass	$\alpha_s(M_Z^2)$
τ -decays	0.1192 ± 0.0018
lattice QCD	0.1188 ± 0.0011
structure functions	0.1156 ± 0.0021
e^+e^- [jets & shps]	0.1169 ± 0.0034
hadron collider	$0.1151^{+0.0028}_{-0.0027}$
ew precision fits	0.1196 ± 0.0030

H1 NNLO jets

0.1156 ± 0.0032